

$$E \frac{d^3 \sigma^{\pi^0}}{dp^3} \text{ [mb / GeV}^2\text{]}$$

■ PHENIX data

GLOBAL **QCD** ANALYSES & PARTON DENSITIES

Rodolfo Sassot
Universidad de Buenos Aires

PHENIX SpinFest 2014

— DSS
- - - KRE
... AKK
▨ scale uncertainty

Fourth Lecture: Nuclear PDFs (nPDFs)

4.0 Prehistory of nPDFs

Experimental and theoretical framework in early '80s

Motivation for studying nuclear effects at high energies

4.1 History of nPDFs

pQCD inspired frameworks, factorization

First nPDFs extractions

4.2 Present status

modern nPDFs: EPS09, nCTEQ, DSSZ

medium modified FFs

4.3 Future of nPDFs

dA (pA) experiments at RHIC (LHC)

Outlook

4.0 Prehistory of nPDFs

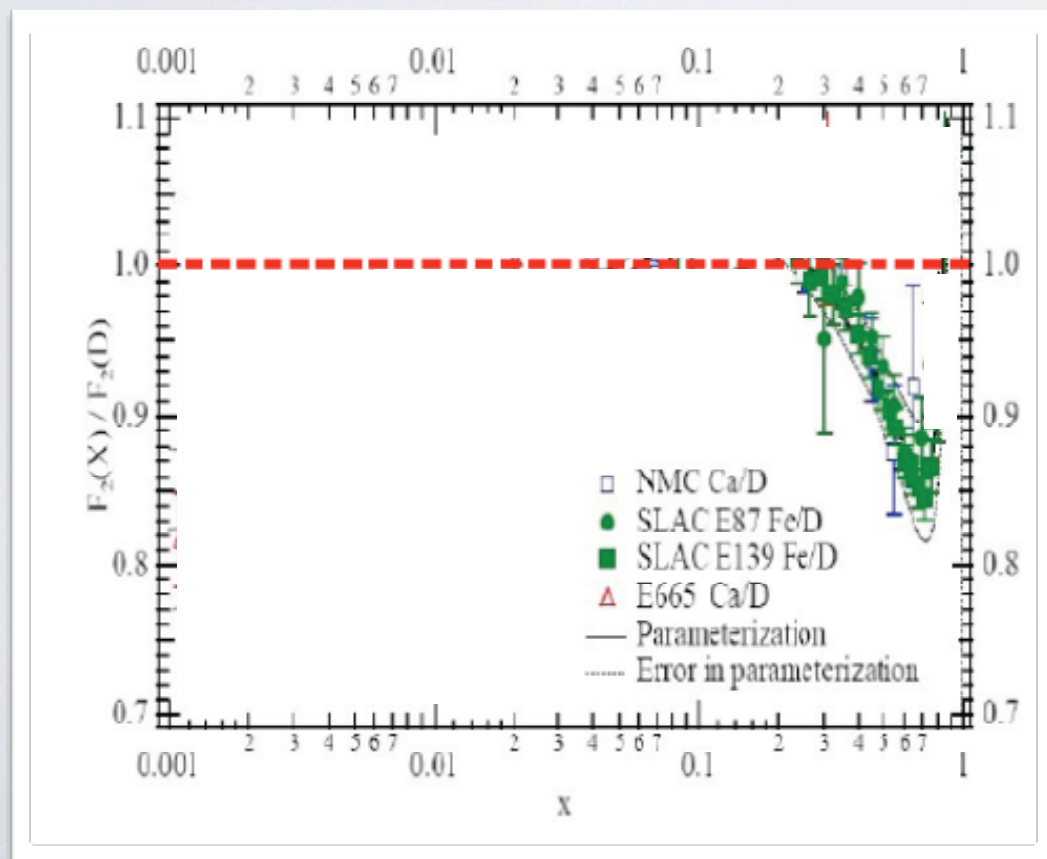
no much interest in partons
in nuclei before 1982:

accurate models for nuclear structure
low energy scales “freeze” QCD dof

DIS “incoherence” hypothesis

$$F_2^A \simeq Z F_2^p + (A - Z) F_2^n$$

European Muon Collaboration (EMC)



$$\frac{\frac{1}{A} F_2^{Fe}}{\frac{1}{2} F_2^d} \neq 1$$

*incoherence?
only nucleons?
free=bounded nucleons?*

interesting in itself but also:

- no neutron targets/beams
- neutrino scattering
- pA, dA baseline for AA
at RHIC and LHC

PDFs

QCD matter

4.0 Prehistory of nPDFs

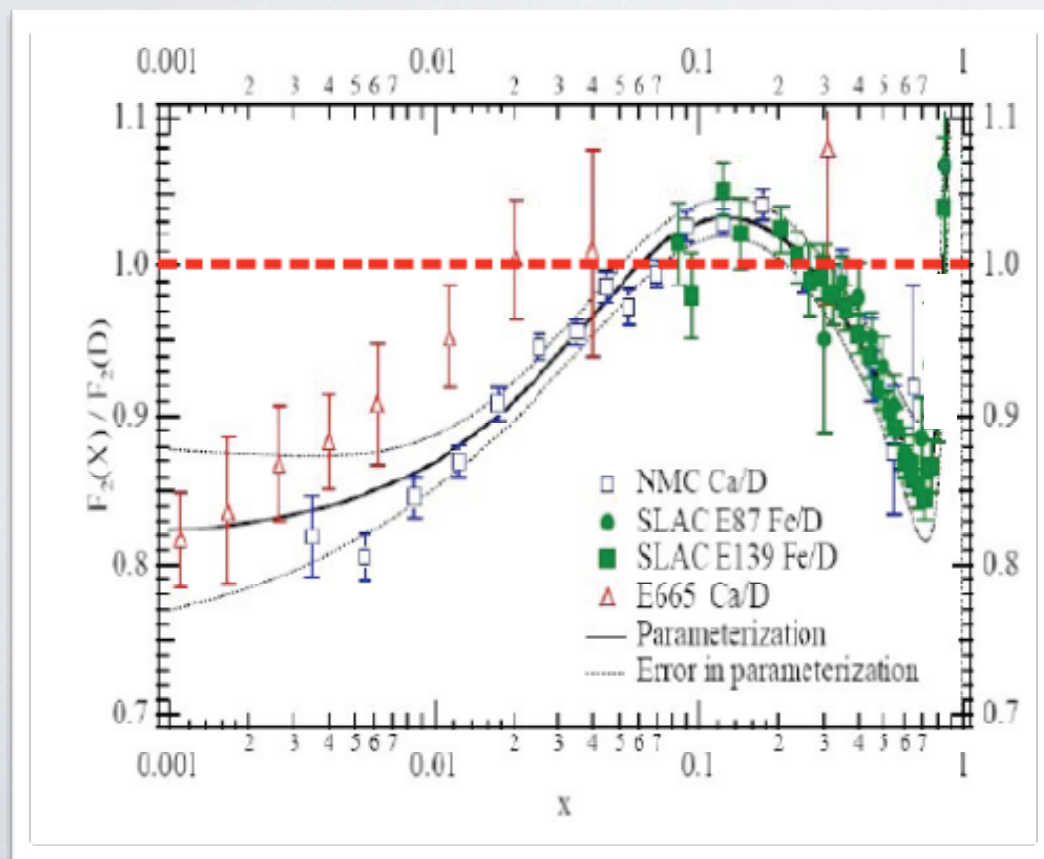
no much interest in partons
in nuclei before 1982:

accurate models for nuclear structure
low energy scales “freeze” QCD dof

DIS “incoherence” hypothesis

$$F_2^A \simeq ZF_2^p + (A - Z)F_2^n$$

European Muon Collaboration (EMC)



$$\frac{\frac{1}{A} F_2^{Fe}}{\frac{1}{2} F_2^d} \neq 1$$

*incoherence?
only nucleons?
free=bounded nucleons?*

interesting in itself but also:

- no neutron targets/beams
- neutrino scattering
- pA, dA baseline for AA
at RHIC and LHC

PDFs

QCD matter

4.0 Prehistory of nPDFs

what is going on?

no universally accepted explanations yet

Fermi motion

- ▶ collective motion of nucleons inside the nucleus
- ▶ enhances “scattering” around & **beyond** (!) kinematic limit for free proton

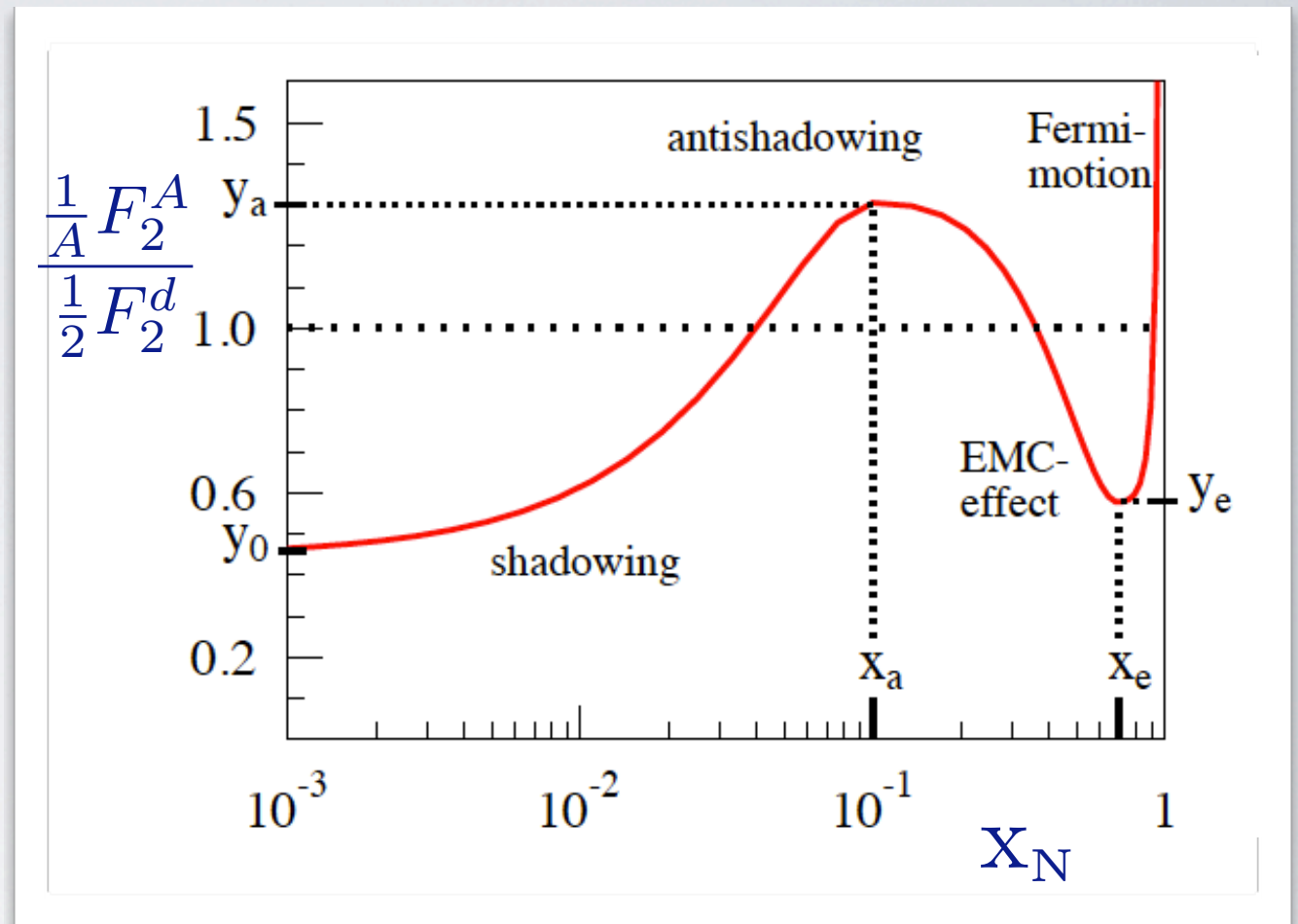
$$\begin{array}{lll} \text{nuclei} & x_A = \frac{Q^2}{2p_A \cdot q} & 0 < x_A < 1 \\ \text{per nucleon} & x_N = \frac{Q^2}{2p_N \cdot q} & 0 < x_N < A \\ p_N = p_A/A & & \end{array}$$

EMC effect

- ▶ ~binding mechanism: if it borrows p_N , works
- ▶ non-nucleonic d.o.f. (pions, multi-quark clusters, ...)
- ▶ many models for bound nucleons

anti-shadowing

- ▶ momentum and baryon number conservation
- ▶ partons (from different nucleons) recombine/fusion

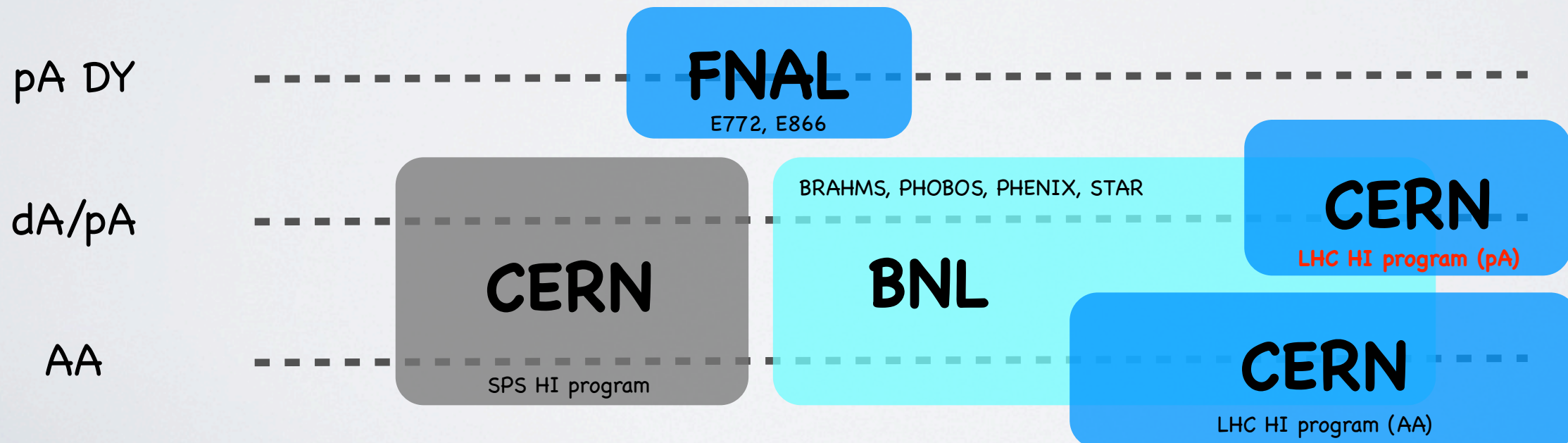
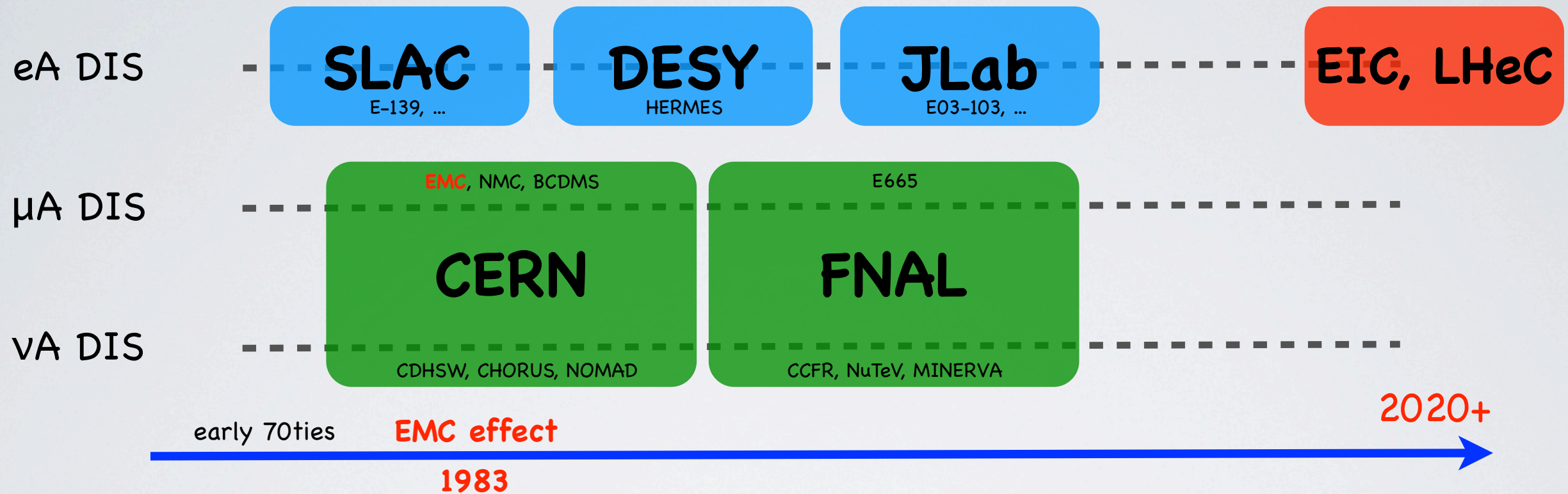


shadowing

- ▶ interpreted as coherent interaction with more than one nucleon, many models
- ▶ effect known in hadron-nucleus total cross sections; optical analogy: surface nucleons shadows inner ones
- ▶ intermediate states: elastic (Glauber) vs inelastic (Gribov)
- ▶ low x ~ parton overlap ~ recombination ~ saturation

4.0 Prehistory of nPDFs

vigorous experimental programs since the early seventies



no eA collider yet

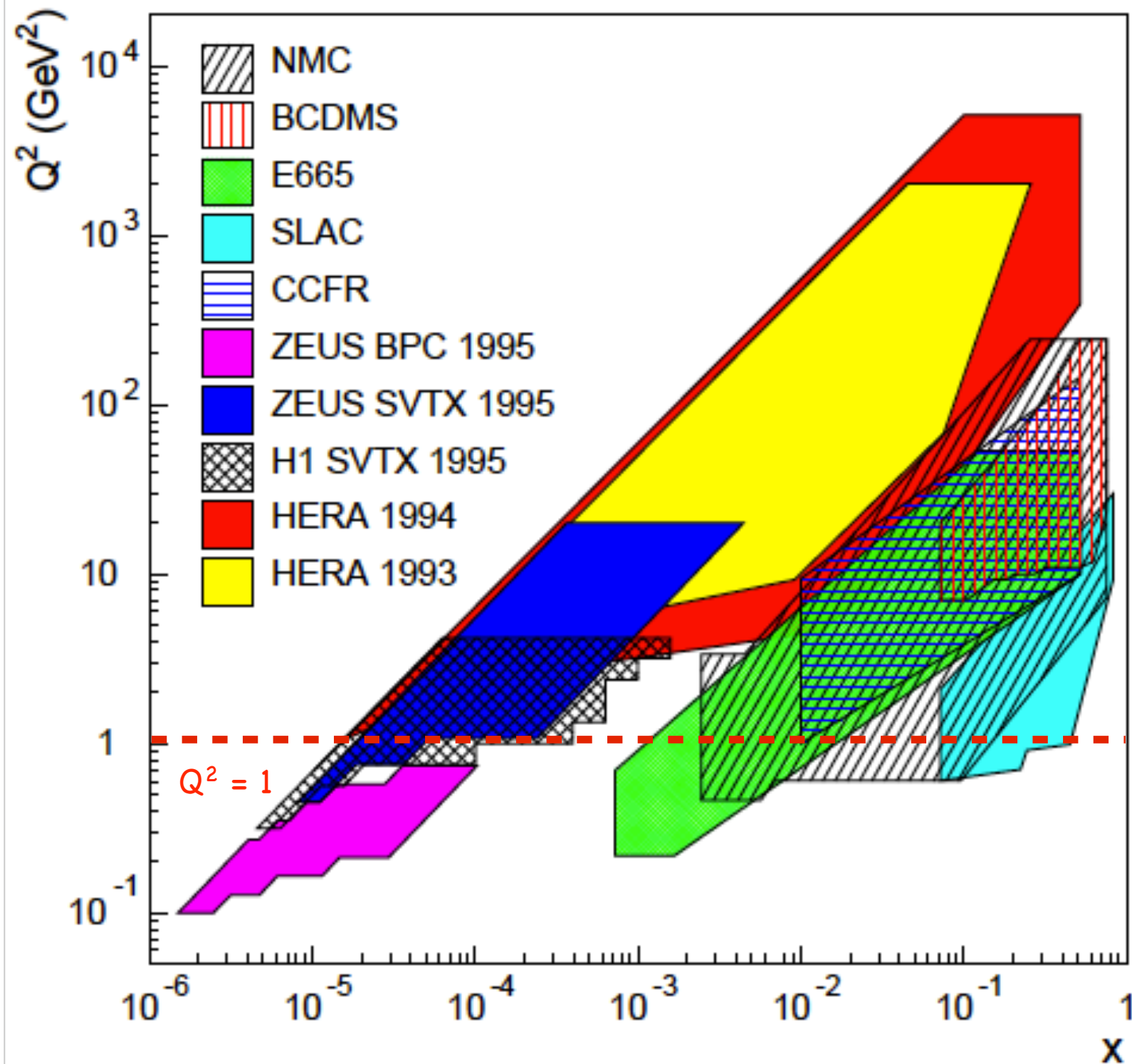
4.0 Prehistory of nPDFs

many different nuclei studied over the years →

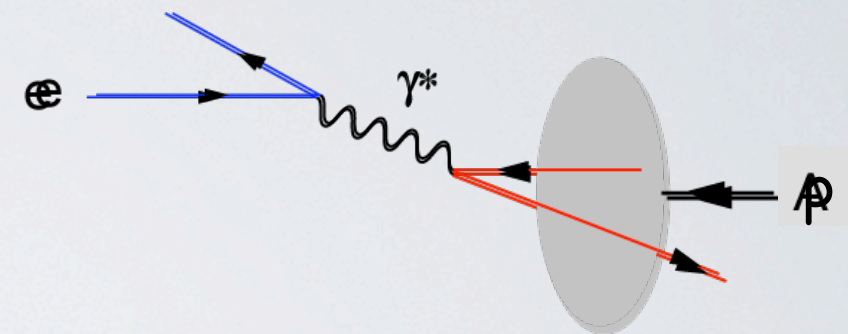
picture of A dependence

[illegible]

4.0 Prehistory of nPDFs



current kinematic coverage
much more limited coverage
for electron-proton DIS
in eA DIS

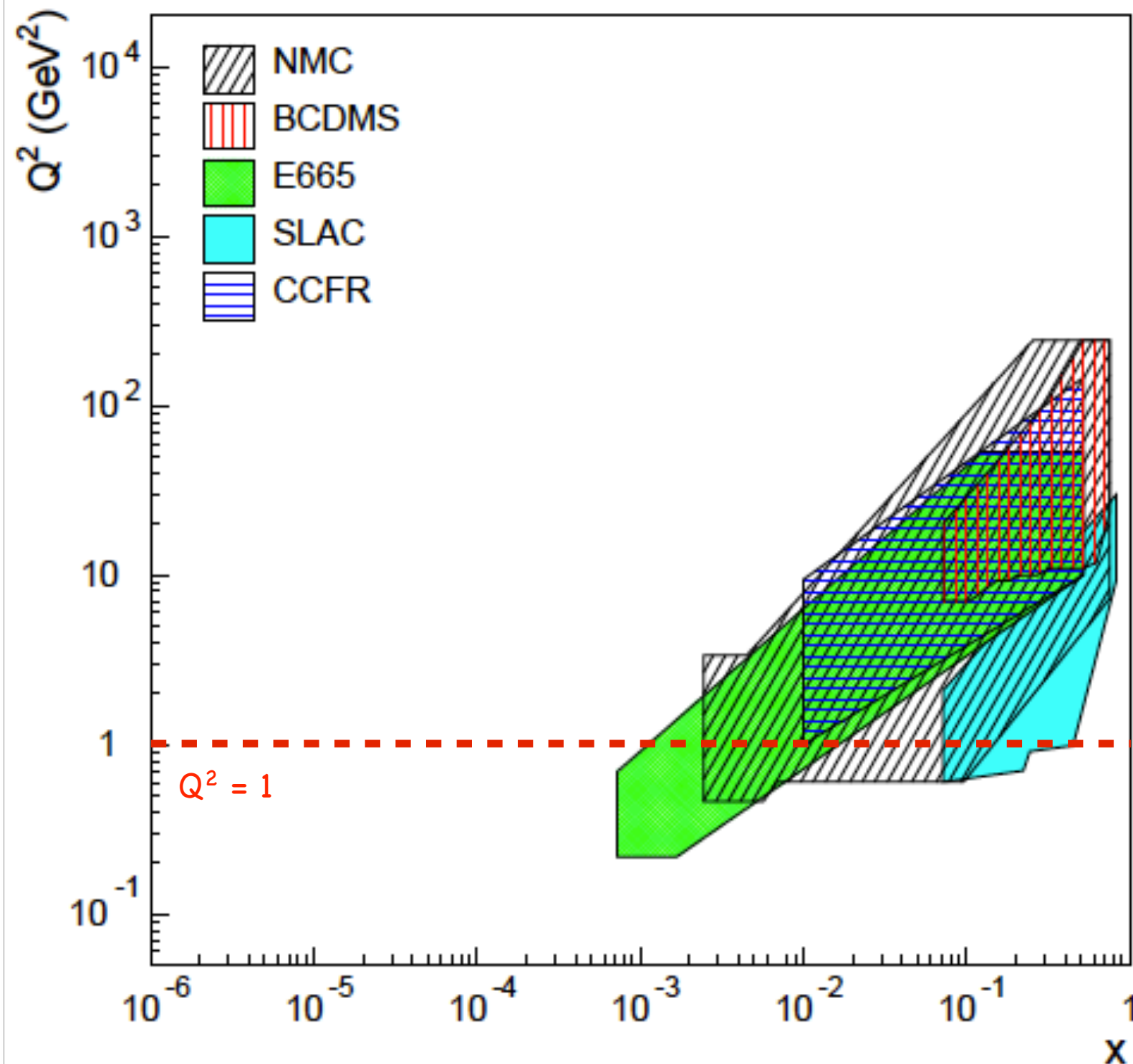


yet the best constraint for nPDFs
determines small- x behaviour
of quarks and gluons in
all analyses of proton PDFs

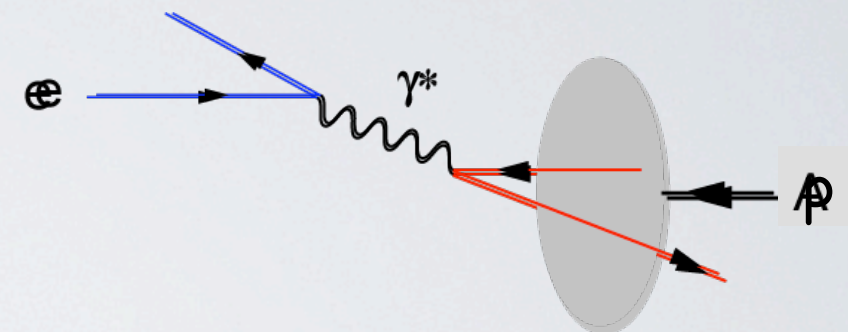
- low x , low Q^2
where saturation is relevant
- high Q^2
to test scale evolution

an electron-ion collider
(EIC, LHeC projects)
is in high demand

4.0 Prehistory of nPDFs



current kinematic coverage
much more limited coverage
for electron-proton DIS
in eA DIS



yet the best constraint for nPDFs
determines small- x behaviour
of quarks and gluons in
all analyses of proton PDFs

- low x , low Q^2
where saturation is relevant
- high Q^2
to test scale evolution

an electron-ion collider
(EIC, LHeC projects)
is in high demand

4.1 History of nPDFs

does a pQCD inspired framework work?

→ factorization between **short** and **long** distances

we calculate

we measure / fit / model

nPDFs definition

$$d\sigma_A = d\hat{\sigma}_i \otimes f_i^A$$

input scale $O(1 \text{ GeV})$

nPDFs parameterization

$$f_i^A(x_N, Q_0^2) = R^A(x_N, Q_0^2) f_i^N(x_N, Q_0^2)$$

free proton

e.g. multiplicative

choose ansatz and determine from data

e.g. from scratch

e.g. convolutions

f_i^A includes all nuclear effects fitted to data

$$d\sigma_N = d\hat{\sigma}_i \otimes f_i^N \text{ PDF}$$

standard partonic cross section

standard nPDFs

standard DGLAP evol. eqs. for nPDFs

"something else"

~medium effects: higher density
non linear
confinement size
 $\Lambda_{\text{QCD}} \dots$

alternative schemes

non-standard evol. eqs. for nPDFs

theorist's favorite playground

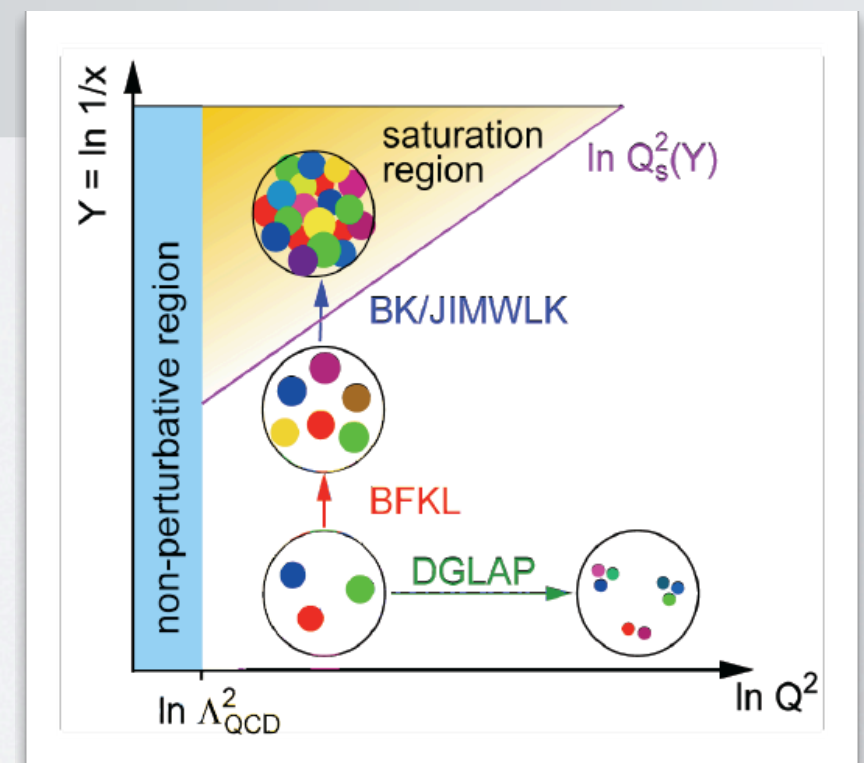
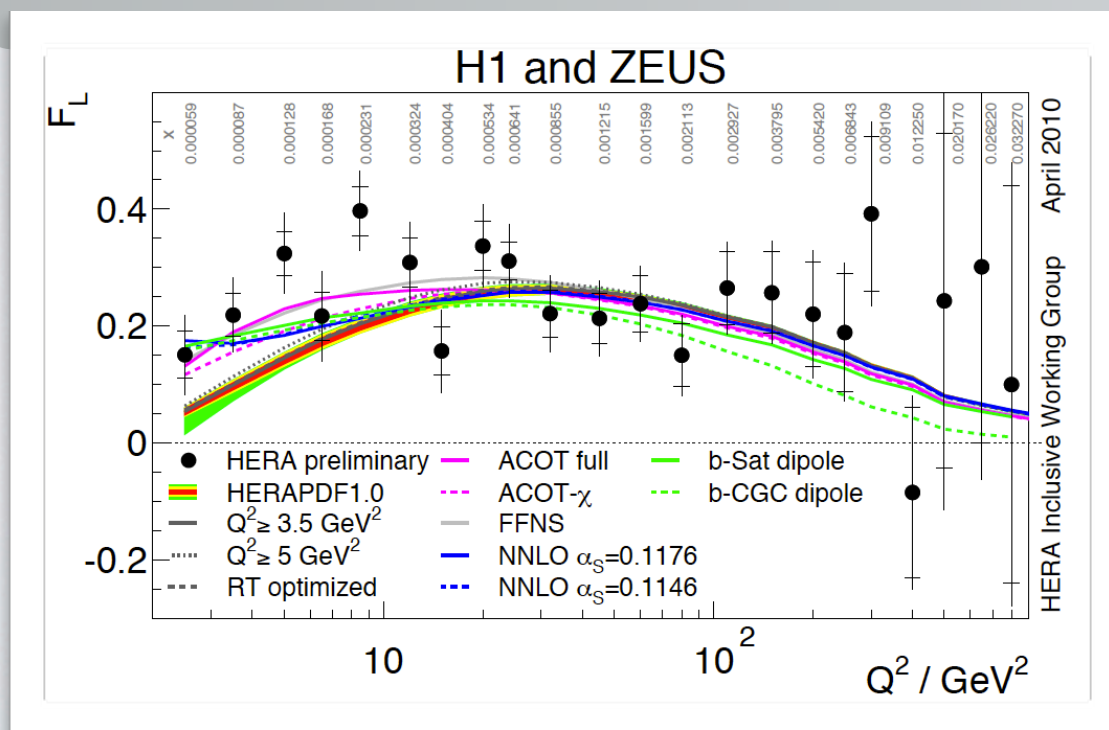
known cross sections
known path to higher orders
universality

other processes?
higher orders?
consistency?

4.1 History of nPDFs

what do we want to learn from nPDFs?

- nPDFs can parametrize nuclear effects with little bias and without assuming certain “mechanisms” to model the observed modifications/effects
link to models of nucleon structure at low scales and proposed nuclear modifications
- a global QCD analysis of many hard probes will reveal tensions due to the assumed framework
factorization and/or DGLAP evolution will eventually break down: where?
- map out kinematic regime where nPDF framework applies and study transition to saturation region

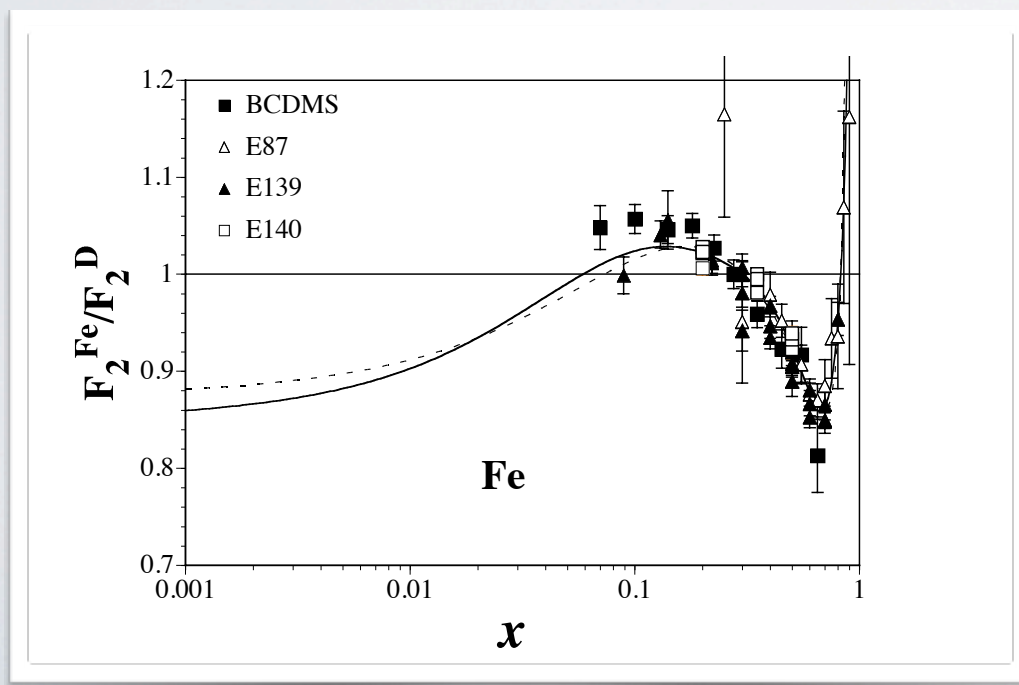
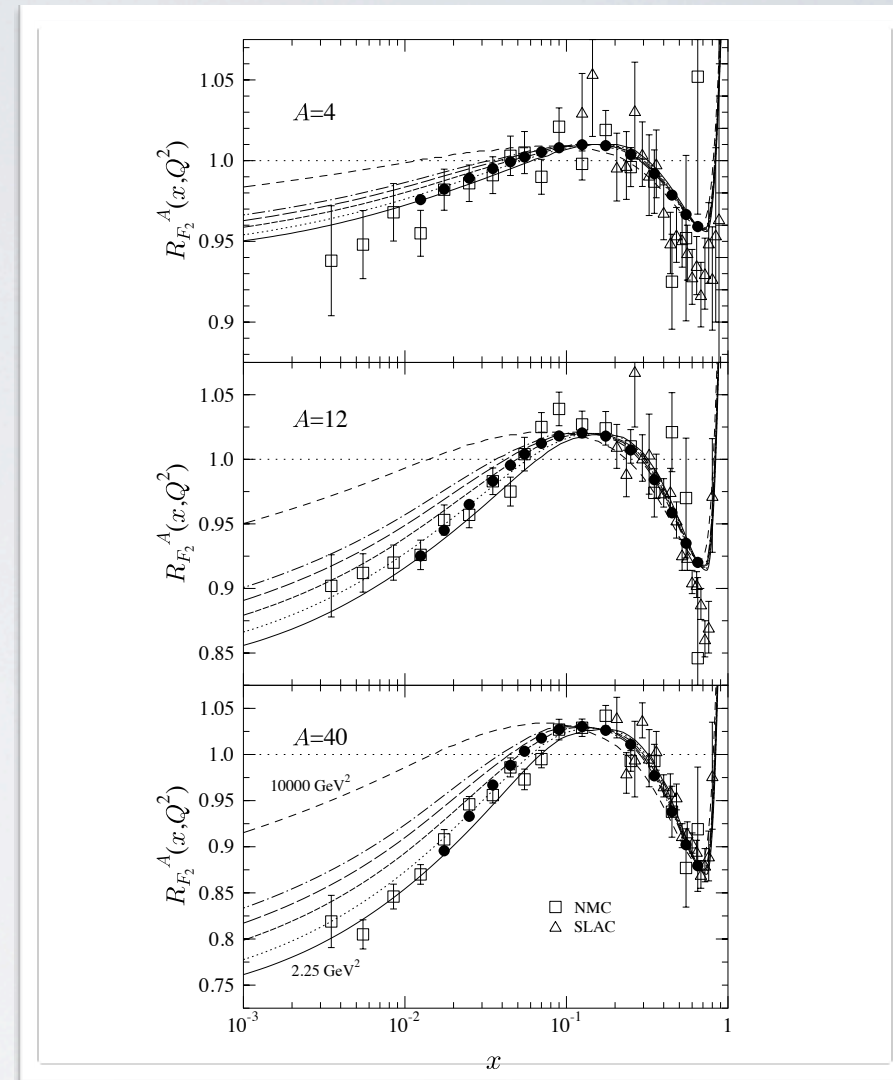


- transition often characterized by “saturation scale” $Q_s(x,A)$
- non-linear effects (recombination) demanded by unitarity
- no unambiguous hints for saturation in ep down to $x = 10^{-5}$

4.1 History of nPDFs

EKS Eskola, Kolhinen, Salgado - hep-ph/9807297
Eskola, Kolhinen, Ruuskanen - hep-ph/9802350

- ▶ **first LO analysis**
- ▶ NMC, E665 DIS and E772 Drell Yan
- ▶ standard multiplicative ansatz
- ▶ no error analysis (no χ^2)



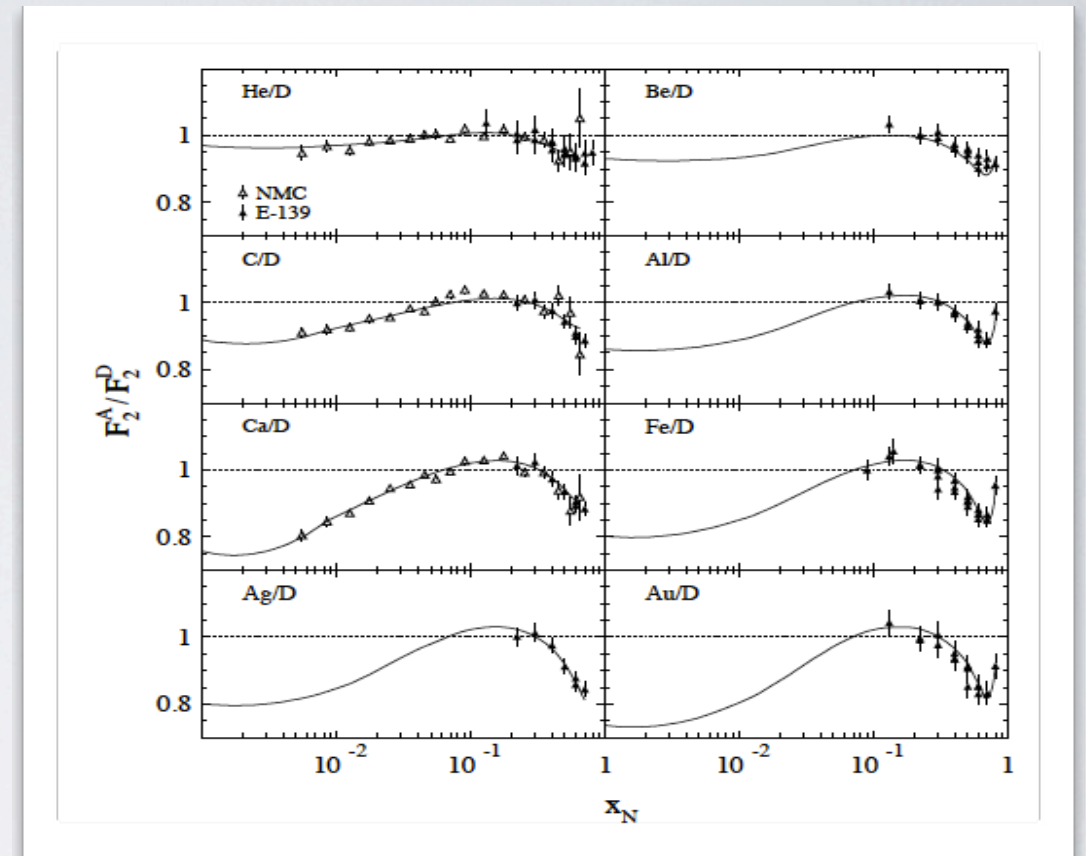
HKN Hirai, Kumano, Nagai - hep-ph/0103208

- ▶ **LO analysis (first χ^2 minimization)**
- ▶ EMC, NMC, SLAC, E665 DIS $\chi^2/\text{d.o.f} = 1.76$
- ▶ standard multiplicative ansatz
- ▶ no error analysis

4.1 History of nPDFs

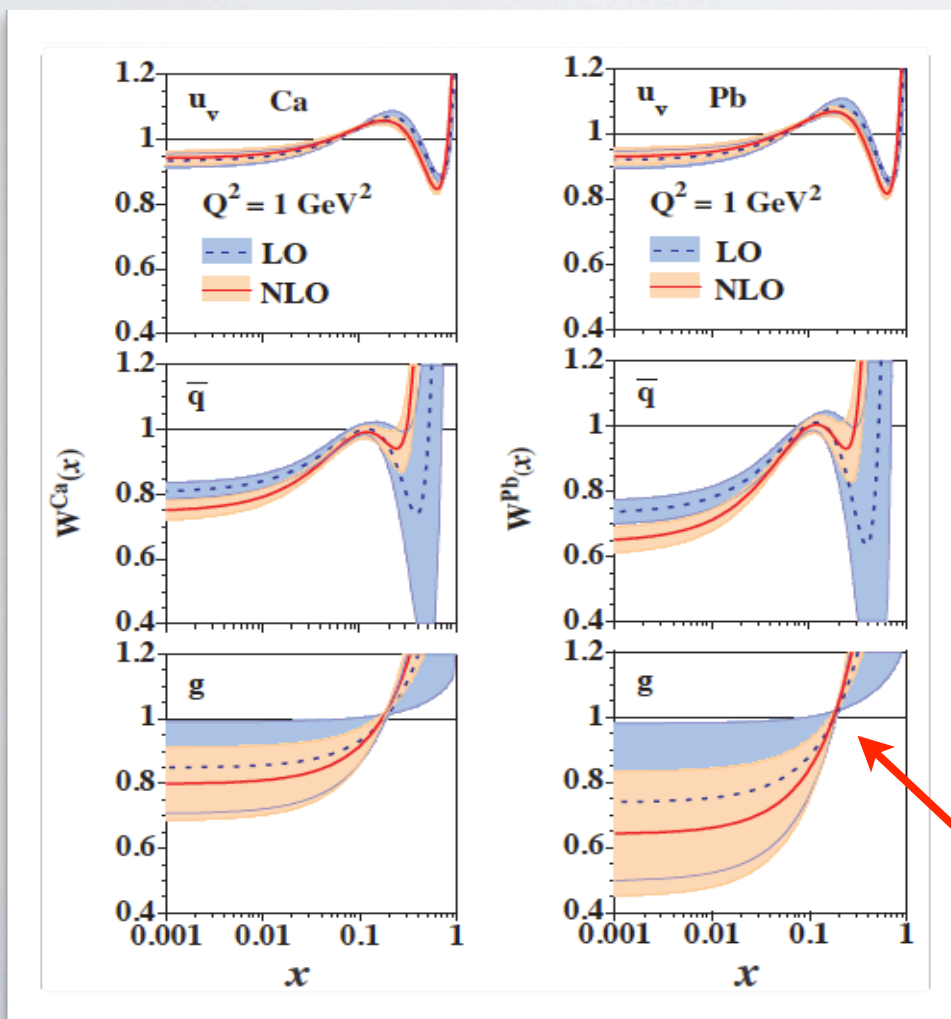
nDS de Florian, R.S. - hep-ph/0311227

- ▶ **first NLO analysis** $\chi^2/\text{d.o.f.} = 0.74$
- ▶ only SLAC & NMC DIS sets and some DY data
- ▶ convolutional approach in Mellin N-space
- ▶ no error analysis



HKN Hirai, Kumano, Nagai - arXiv:0709.3038

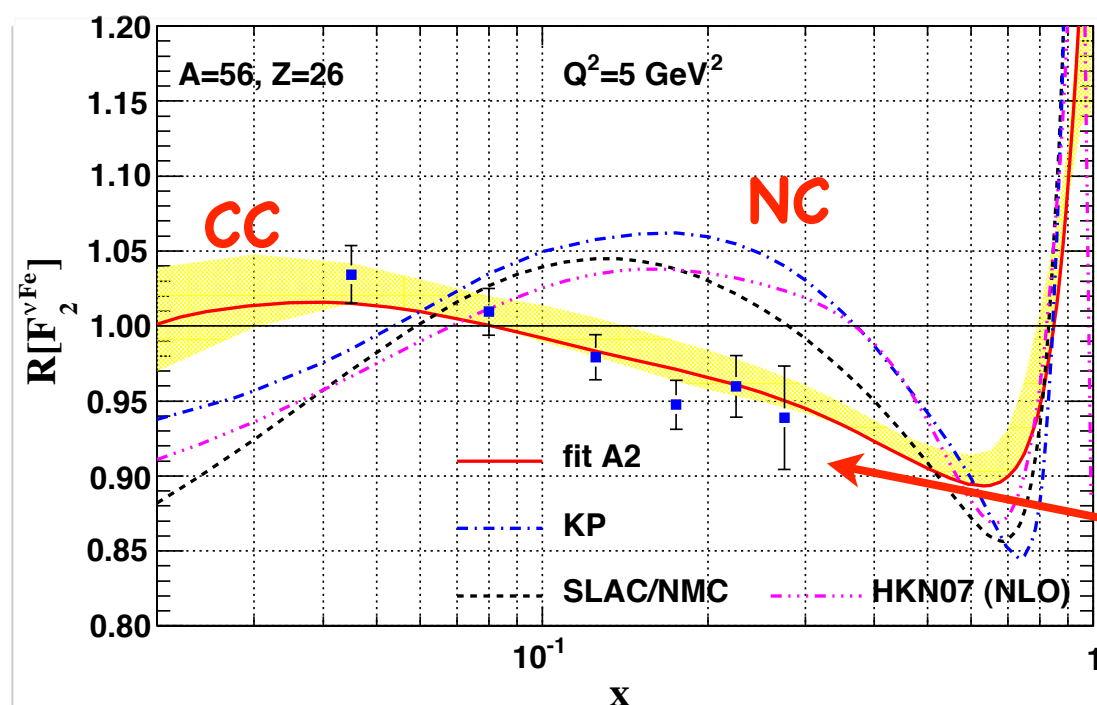
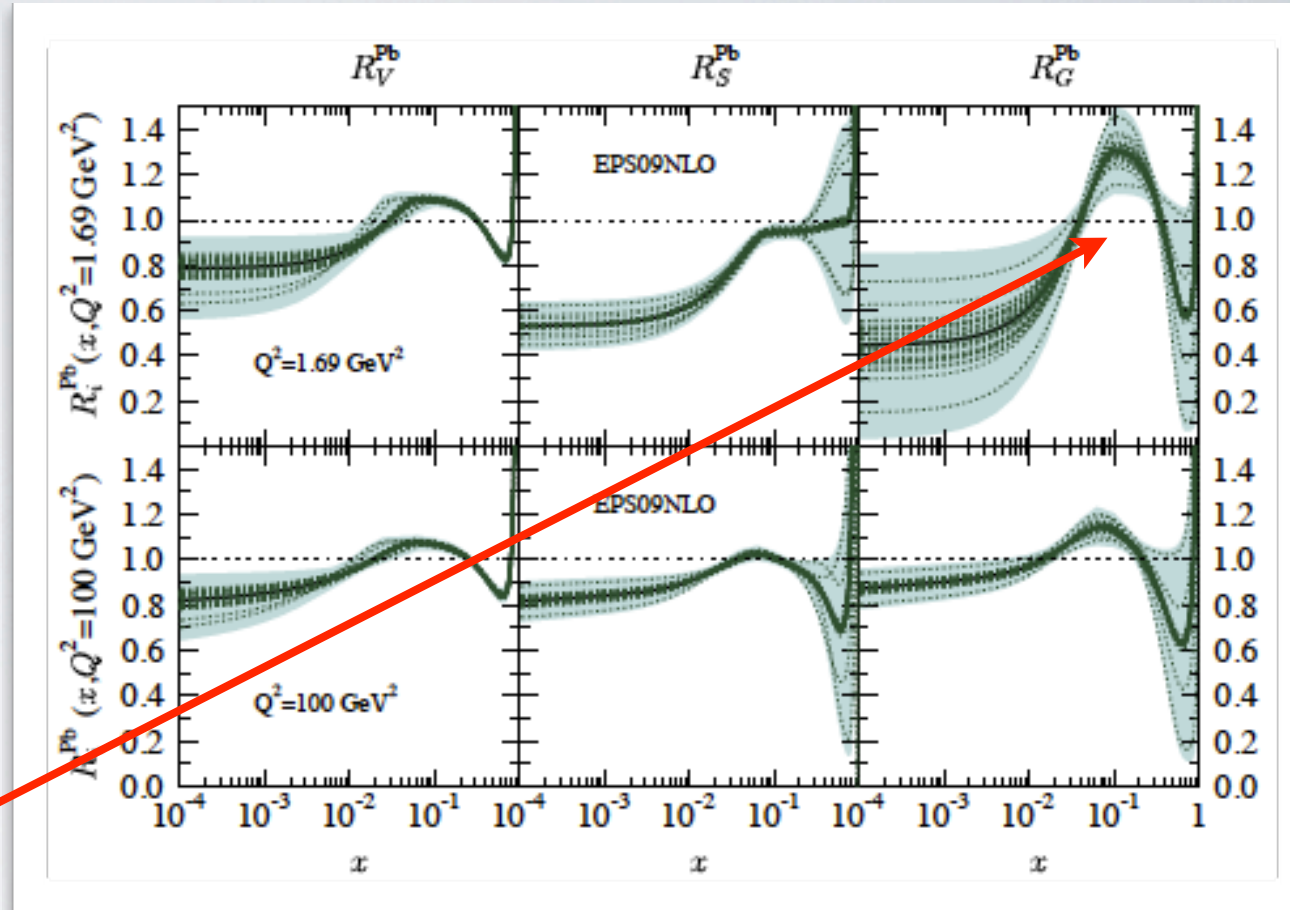
- ▶ LO and NLO analyses $\chi^2/\text{d.o.f.} = 1.2$
- ▶ standard DIS and DY data sets
- ▶ standard multiplicative ansatz
- ▶ **first error analysis** (Hessian method)
- ▶ rather "unusual" gluon distribution at large x



4.2 Present status

EPS Eskola, Paukkunen, Salgado - 0902.4154

- ▶ NLO analysis $\chi^2/\text{d.o.f.} = 0.8$
- ▶ usual DIS & DY data
- ▶ **RHIC dAu data to constrain gluon better**
- ▶ complicated piecewise multipl. ansatz
- ▶ Hessian error analysis
- ▶ huge anti-shadowing/EMC effect for gluon



nCTEQ Keppel, Kovarik, ... - 0907.2357

- ▶ direct ansatz a la CTEQ
- ▶ DIS & DY plus **CC neutrino DIS data**
- ▶ find tension between NC and CC DIS data

breakdown of factorization

4.2 Present status: DSSZ

de Florian, R.S., Stratmann, Zurita - 1112.6324

why do we need yet another set of nPDFs ?

▶ no truly global analysis yet

↳ *include charged lepton DIS, Drell-Yan, CC neutrino DIS, and RHIC dAu data*

▶ use up-to-date proton PDFs as reference set

↳ *many different sets to choose from - take MSTW*

Martin, Stirling, Thorne, Watt - arXiv:0901.0002

▶ improve on the treatment of heavy flavors

↳ *e.g. NLO massive Wilson cross sections for CC DIS*

Blumlein, Hasselhuhn, Kovacikova, Moch - arXiv:1104.3449

▶ provide some estimate of nPDF uncertainties

main questions to address

- *do we really see a tension between charged lepton and neutrino DIS data?*
- *do RHIC dAu data imply strong modifications of the nuclear gluon distribution?*

DSSZ: x-dependence

- ▶ use multiplicative nuclear modification factor $f_i^A(x, Q_0) = R_i^A(x, Q_0) \times f_i^P(x, Q_0)$
- ▶ initial scale $Q_0 = 1 \text{ GeV}$, NLO DGLAP evolution to all other scales $Q > Q_0$

- ▶ parametrize both valence distributions as needs to be flexible enough to accommodate (anti-)shadowing, EMC effect, Fermi motion

$$R_v^A(x, Q_0) = \epsilon_1 x^{\alpha_v} (1-x)^{\beta_1} \times [1 + \epsilon_2 (1-x)^{\beta_2}] \times [1 + a_v (1-x)^{\beta_3}]$$

- ▶ data do not allow to discriminate different sea quark flavors (tried in analysis)

$$R_s^A(x, Q_0) = R_v^A(x, Q_0) \frac{\epsilon_s}{\epsilon_1} \frac{1 + a_s x^{\alpha_s}}{a_s + 1}$$

- ▶ need another modification factor for gluons

$$R_g^A(x, Q_0) = R_v^A(x, Q_0) \frac{\epsilon_g}{\epsilon_1} \frac{1 + a_g x^{\alpha_g}}{a_g + 1}$$

- ▶ heavy quarks generated radiatively: no parameters

- ▶ 3 parameters constrained by charge & momentum conservation

also, fit unchanged if
 $\epsilon_g = \epsilon_s$

total of 9 parameters per nucleus

$$\xi \in \{\alpha_v, \alpha_s, \alpha_g, \beta_1, \beta_2, \beta_3, a_v, a_s, a_g\}$$

quality of the fit unchanged
by relating $R_{s,g}$ to common R_v

but need different
normalization and small- x behavior

DSSZ: A-dependence

total of 9 parameters per nucleus

$$\xi \in \{\alpha_v, \alpha_s, \alpha_g, \beta_1, \beta_2, \beta_3, a_v, a_s, a_g\}$$

► A dependence implemented as

$$\xi = \gamma_\xi + \lambda_\xi A^{\delta_\xi}$$

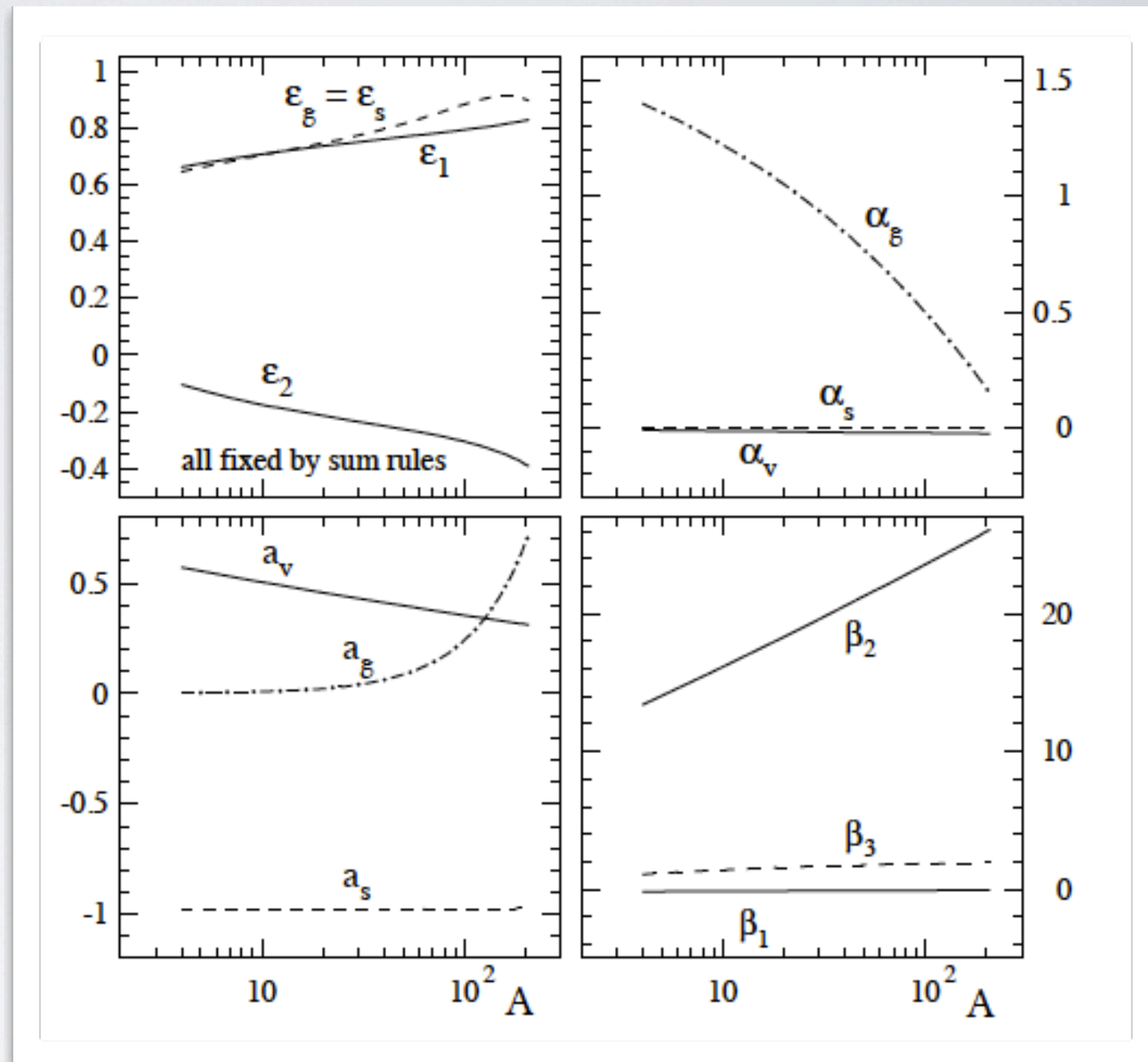
► fit does not fix all parameters, assume

$$\delta_{a_g} = \delta_{a_s} \quad \delta_{\alpha_g} = \delta_{\alpha_s}$$

**25 free parameters
in total**

parameter	γ	λ	δ
α_v	-0.256	0.252	-0.017
α_s	0.001	-6.89×10^{-4}	0.286
α_g	1.994	-0.401	0.286
β_1	-5.564	5.36	0.0042
β_2	-59.62	69.01	0.0407
β_3	2.099	-1.878	-0.436
a_v	-0.622	1.302	-0.062
a_s	-0.980	2.33×10^{-6}	1.505
a_g	0.0018	2.35×10^{-4}	1.505

A dependence of fit parameters



optimum NLO parameters
at the input scale

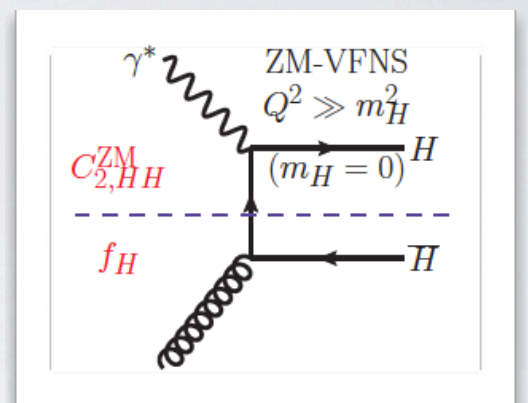
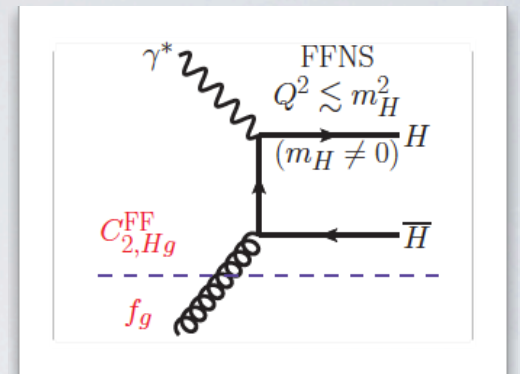
DSSZ: heavy flavors

different ways to treat heavy quarks:

- $Q \not\gg m_Q$ \uparrow **fixed flavor number scheme (FFNS)**
 HQ only produced extrinsically via, e.g., photon-gluon fusion
- $Q \gg m_Q$ \uparrow **zero mass variable flavor number scheme (ZM-VFNS)**
 massless HQ above "threshold" $Q = m_Q$; active in evolution
- $Q \gg m_Q$ \uparrow **general mass variable flavor number scheme (GM-VFNS)**
 attempt to match massless and massive theories;
 needs some matching & interpolating coefficients
 details vary in global fits (many prescriptions)

not a priori clear if/where it matters

e.g. HERA data described well in both FFNS and GM-VFNS; ZM-VFNS clearly inadequate



charm production in CC DIS is of special interest

$$W^+ s' \rightarrow c \quad s' \equiv |V_{cs}|^2 s + |V_{cd}|^2 d$$

➡ *used to extract strangeness from CC neutrino data in PDF fits*

need to control nuclear corrections for Fe and Pb targets

➡ *we adopt the GM-VFNS as defined in the free proton PDFs of MSTW*

positive impact on quality of our fit to CC DIS data: 26% gain in χ^2

DSSZ: data sets & χ^2

- ▶ optimum parameters determined from standard chi-squared optimization

relative normalization or
artificial weights for certain data sets

not needed/used
in DSSZ analysis

$$\chi^2 \equiv \sum_i \omega_i \frac{(\text{d}\sigma_i^{\text{exp}} - \text{d}\sigma_i^{\text{th}})^2}{\Delta_i^2}$$

uncertainty for each point

DSSZ: add sys + stat in quadrature [+ theor. unc.]

total

$\chi^2 : 1544.7/1579\text{pts.}$
 $\chi^2/\text{d.o.f} : 0.994$

measurement	collaboration	# points	χ^2	
F_2^{He}/F_2^D	NMC	17	18.18	NC DIS 897.5/894
	E139	18	2.71	
F_2^{Li}/F_2^D	NMC	17	17.35	
F_2^{Li}/F_2^D Q^2 dep.	NMC	179	197.36	
F_2^{Be}/F_2^D	E139	17	44.17	
F_2^C/F_2^D	NMC	17	27.85	
	E139	7	9.66	
	EMC	9	6.41	
F_2^C/F_2^D Q^2 dep.	NMC	191	201.63	
F_2^{Al}/F_2^D	E139	17	13.22	
F_2^{Ca}/F_2^D	NMC	16	18.60	
	E139	7	12.13	
F_2^{Cu}/F_2^D	EMC	19	18.62	
F_2^{Fe}/F_2^D	E139	23	34.95	
F_2^{Ag}/F_2^D	E139	7	9.71	
F_2^{Sn}/F_2^D	EMC	8	16.59	
F_2^{Au}/F_2^D	E139	18	10.46	
F_2^C/F_2^{Li}	NMC	24	33.17	
F_2^{Ca}/F_2^{Li}	NMC	24	25.31	
F_2^{Be}/F_2^C	NMC	15	11.76	
F_2^{Al}/F_2^C	NMC	15	6.93	CC DIS 488.2/532
F_2^{Ca}/F_2^C	NMC	15	7.71	
F_2^{Ca}/F_2^C	NMC	24	26.09	
F_2^{Fe}/F_2^C	NMC	15	10.38	
F_2^{Sn}/F_2^C	NMC	15	4.69	
F_2^{Sn}/F_2^C Q^2 dep.	NMC	145	102.31	
F_2^{Pb}/F_2^C	NMC	15	9.57	
F_2^{vFe}	NuTeV	78	109.65	Drell Yan 90.7/92
F_3^{vFe}	NuTeV	75	79.78	
F_2^{vFe}	CDHSW	120	108.20	
F_3^{vFe}	CDHSW	133	90.57	
F_2^{vPb}	CHORUS	63	20.42	
F_3^{vPb}	CHORUS	63	79.58	
$d\sigma_{DY}^C/d\sigma_{DY}^D$	E772	9	9.87	
$d\sigma_{DY}^{Ca}/d\sigma_{DY}^D$	E772	9	5.38	dAu->piX 68.3/61
$d\sigma_{DY}^{Fe}/d\sigma_{DY}^D$	E772	9	9.77	
$d\sigma_{DY}^W/d\sigma_{DY}^D$	E772	9	19.29	
$d\sigma_{DY}^{Fe}/d\sigma_{DY}^{Be}$	E866	28	20.34	
$d\sigma_{DY}^W/d\sigma_{DY}^{Be}$	E866	28	26.07	
$d\sigma_{\pi^0}^{dAu}/d\sigma_{\pi^0}^{pp}$	PHENIX	20	27.71	
$d\sigma_{\pi^0}^{dAu}/d\sigma_{\pi^0}^{pp}$	STAR	11	3.92	
$d\sigma_{\pi^\pm}^{dAu}/d\sigma_{\pi^\pm}^{pp}$	STAR	30	36.63	
Total		1579	1544.70	

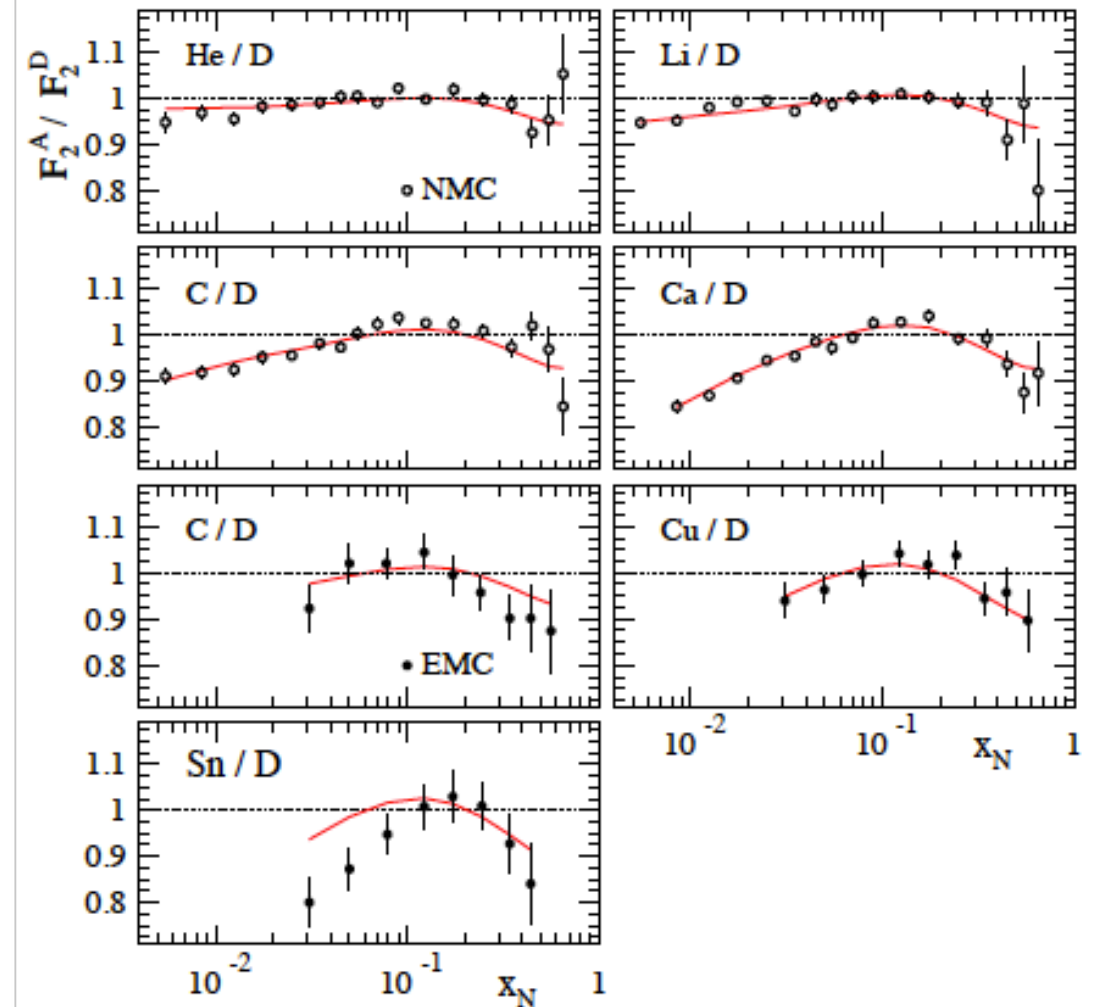
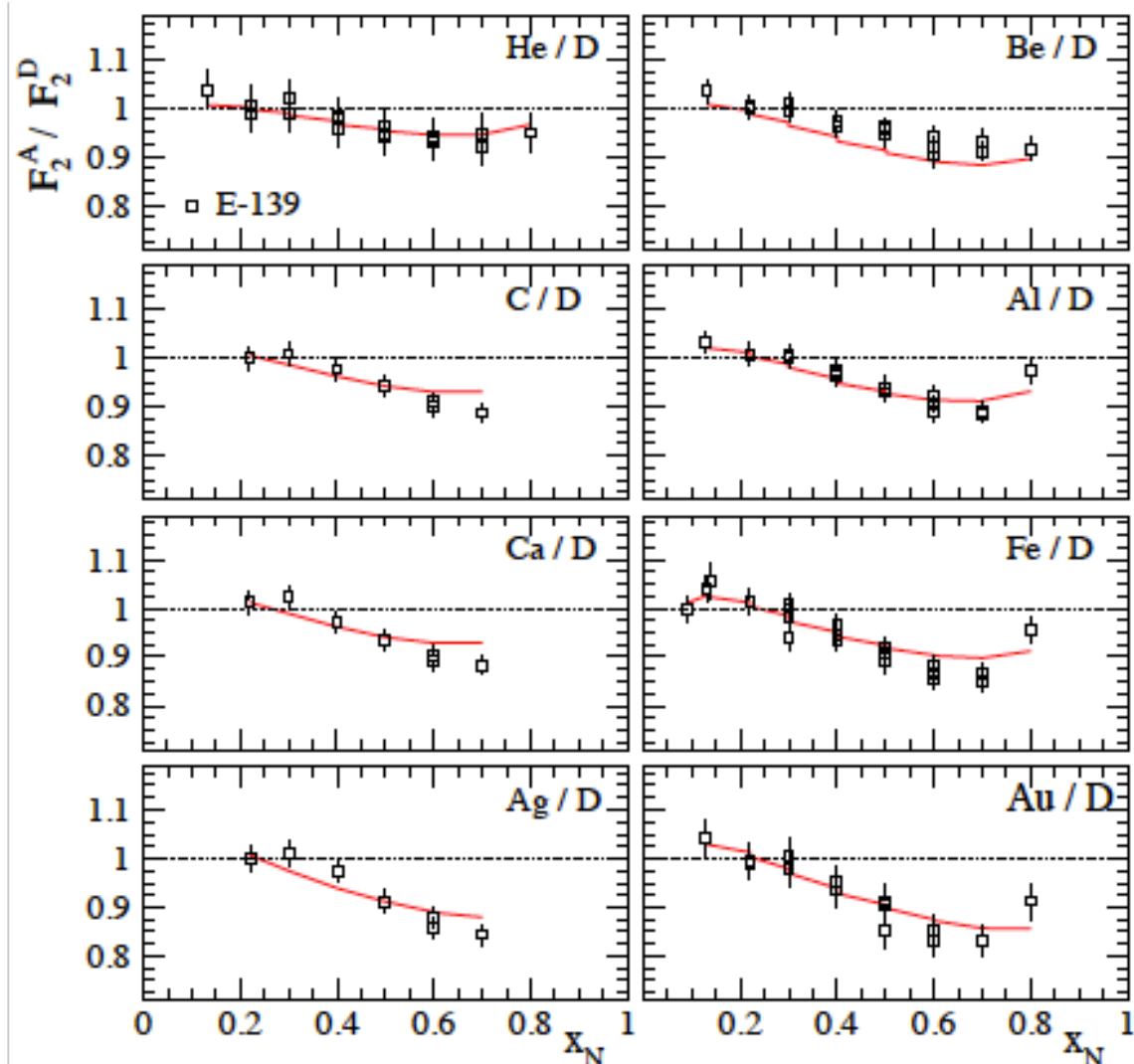
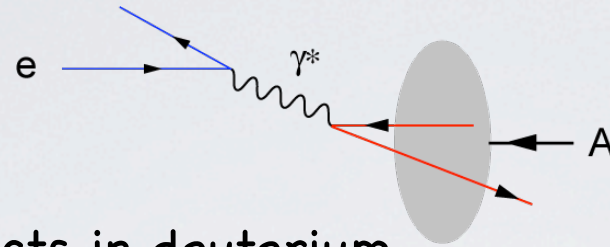
DSSZ: charged lepton DIS data

fit all “classic” EMC, NMC, and E-139 DIS data

► impose cut $Q^2 > 1 \text{ GeV}^2$

► $\chi^2 = 857.5/894\text{pts.}$

► neglect, as usual, nuclear effects in deuterium found to be small in Hirai, Kumano, Nagai



recall

main constraint
from DIS data $0.01 \lesssim x \lesssim 0.8$

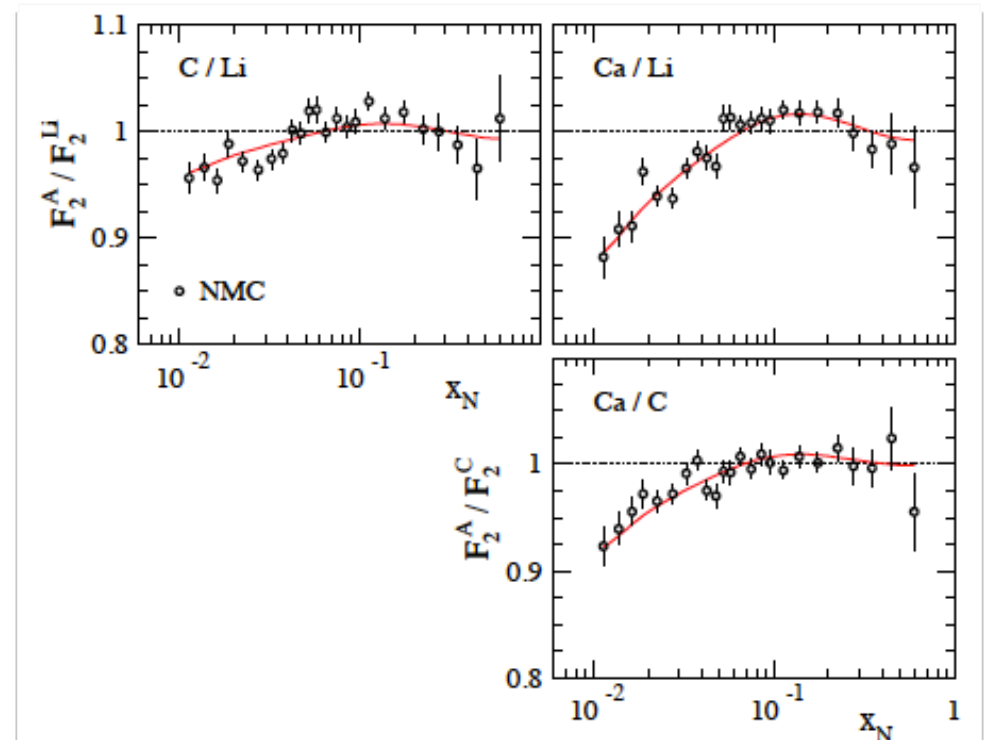
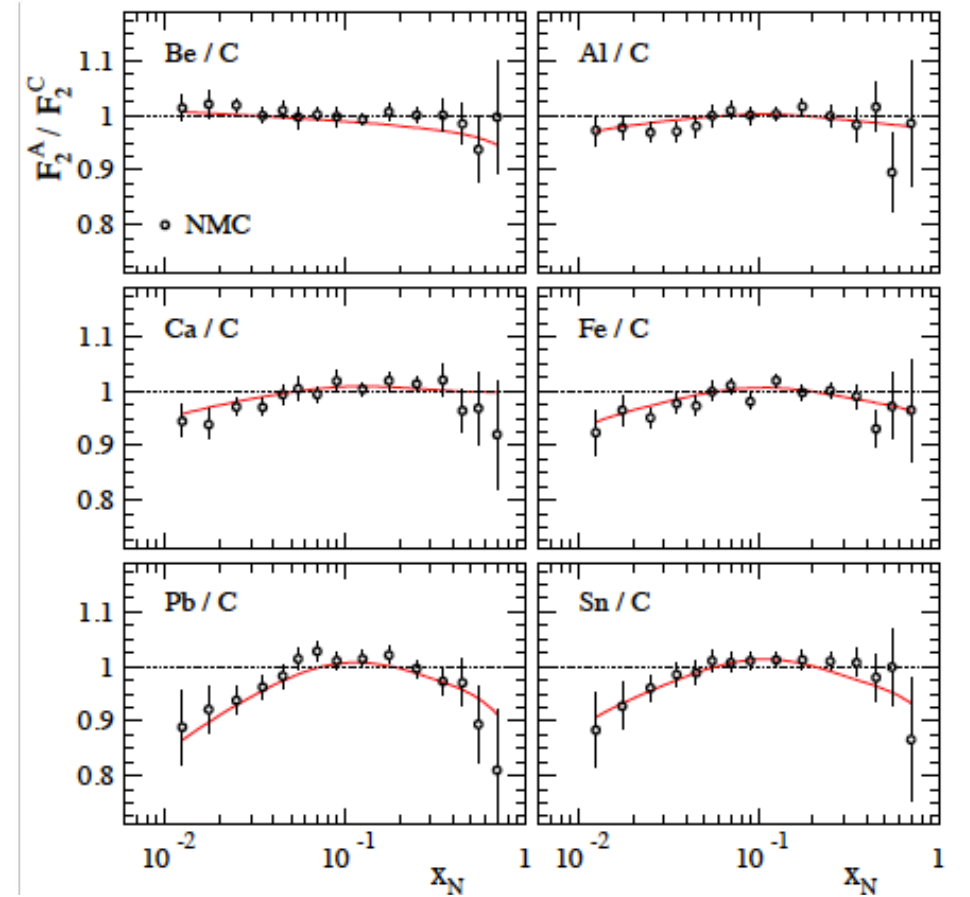
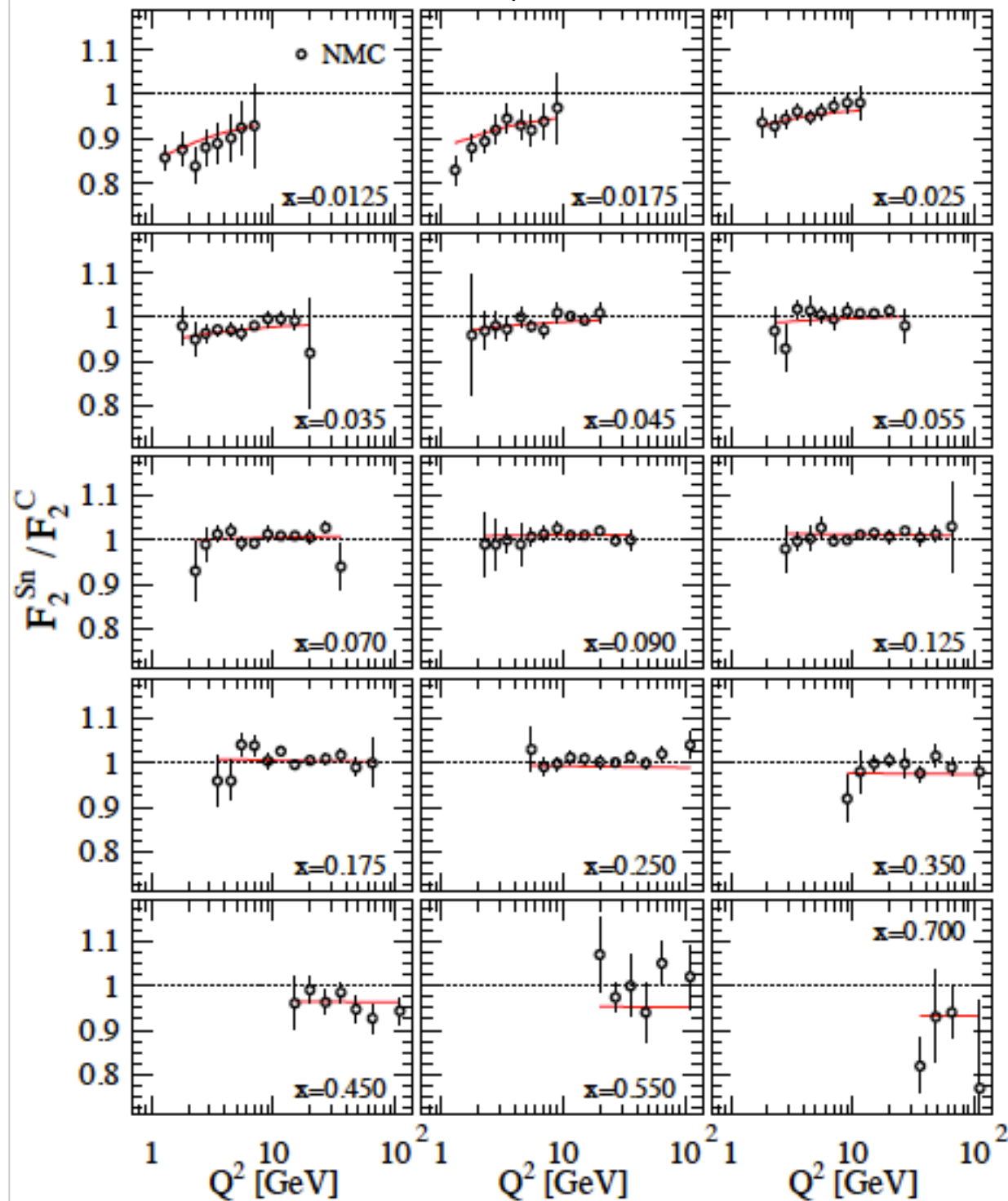
$$F_2^A(N) = x \sum_q e_q^2 \left[(q^A(N) + \bar{q}^A(N)) \left(1 + \frac{\alpha_s}{2\pi} C_2^q(N) \right) + \frac{\alpha_s}{2\pi} C_2^g(N) g^A(N) \right]$$

weak indirect constraint
from scale evolution

DSSZ: charged lepton DIS data

there is more ... no surprises though

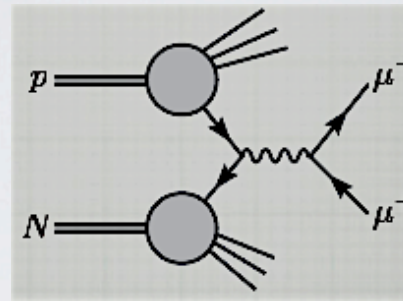
similarly for $F_2^{\text{Li}}/F_2^{\text{D}}$ and $F_2^{\text{C}}/F_2^{\text{D}}$



DSSZ: Drell-Yan dimuon data

fit all **E772** and **E866** DY pA data

- ▶ di-muons have inv. mass $M > 4$ GeV (sets scale)
- ▶ $\chi^2 = 90.7/92\text{pts.}$



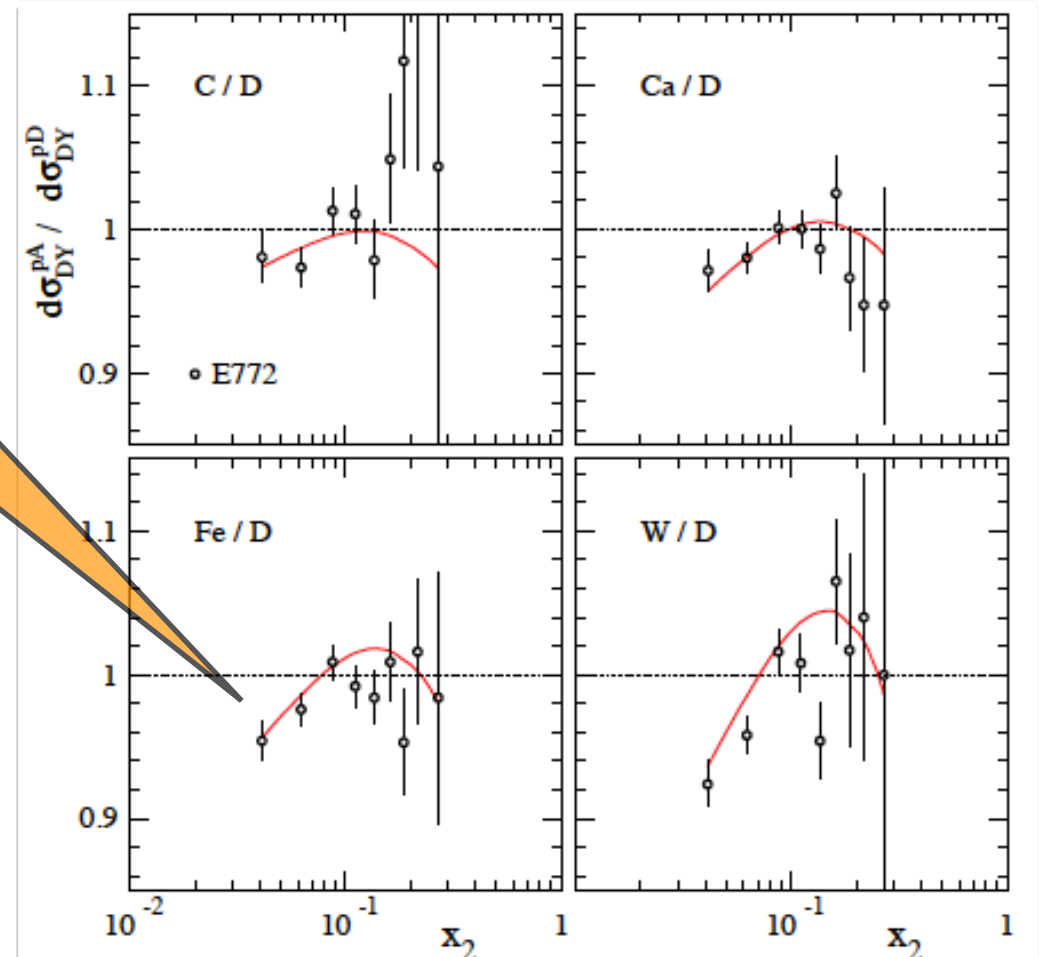
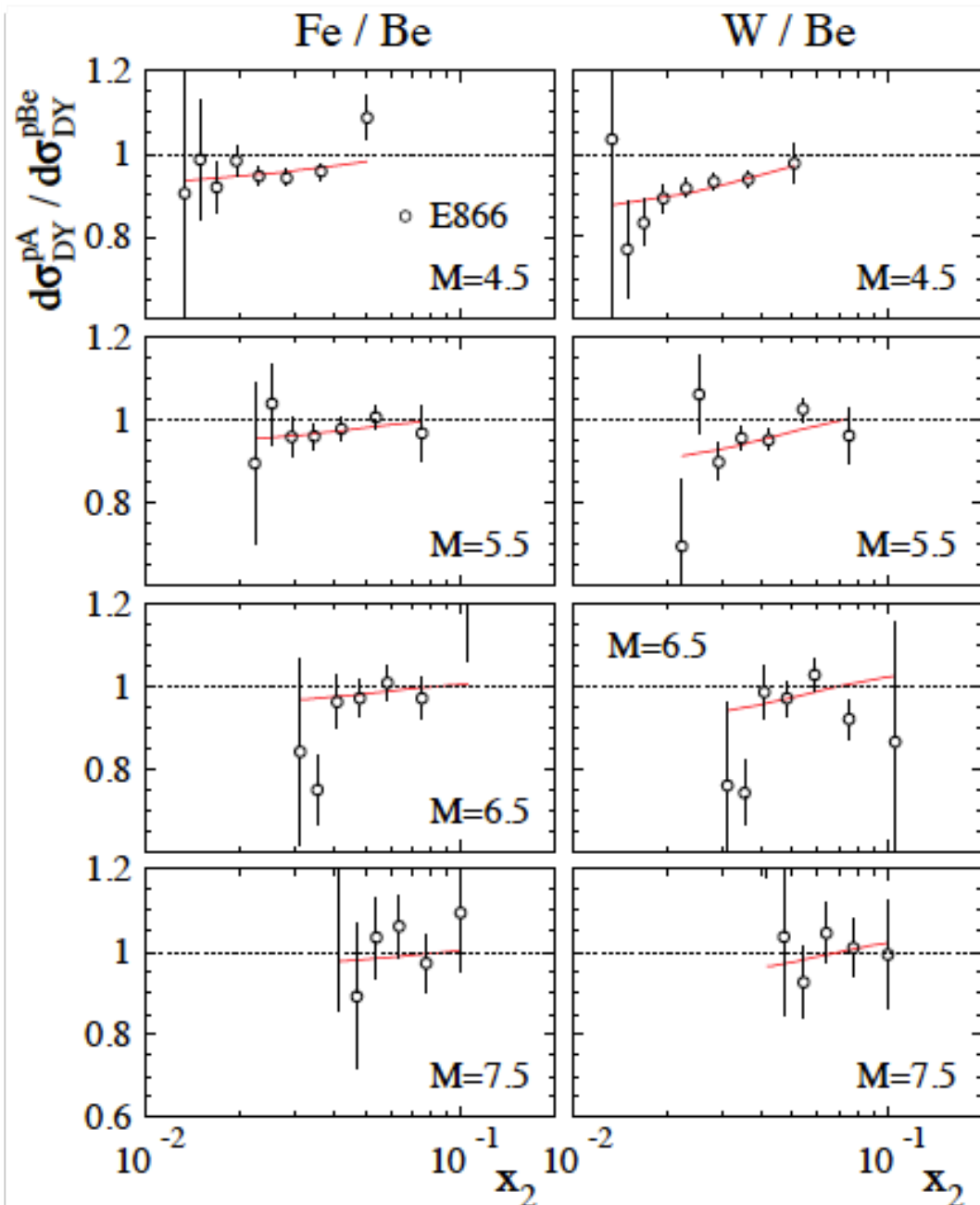
$$x_{1,2} = \sqrt{M^2/s} e^{\pm y}$$

DY data mainly help to
disentangle val/sea quarks
gluons through evolution

$$\frac{d^2\sigma}{dMdy} = \frac{4\pi\alpha^2}{9M^3} \sum_{ij} \int dx_1 dx_2 f_i^p(x_1) f_j^A(x_2) \frac{d\hat{\sigma}_{ij}}{dMdy}$$

$x_2 \in [0.01, 0.2]$

"evidence"
for shadowing
of sea quarks



DSSZ: cc neutrino DIS data

fit **CDHSW**, **NuTeV**, and **CHORUS** str. fct. data

substantial interest:

- **nCTEQ** claim of “**factorization breaking**” for nPDFs
- neutrino data are a vital constraint on strangeness (and help to separate quark flavors) in proton PDF fits

$$\frac{d^2\sigma^{\nu A, \bar{\nu} A}}{dx dy} \simeq xy^2 F_1^{\nu A, \bar{\nu} A} + (1-y) F_2^{\nu A, \bar{\nu} A} \pm xy(1 - \frac{y}{2}) F_3^{\nu A, \bar{\nu} A}$$

- CC DIS data probe different combinations of up-/down-type quarks than charged-lepton DIS
- neutrino and antineutrino beams probe 4 different structure functions

$$F_2^{\nu A}(x_N) \simeq x_N [\bar{u}^A + \bar{c}^A + d^A + s^A](x_N)$$

$$F_2^{\bar{\nu} A}(x_N) \simeq x_N [u^A + c^A + \bar{d}^A + \bar{s}^A](x_N)$$

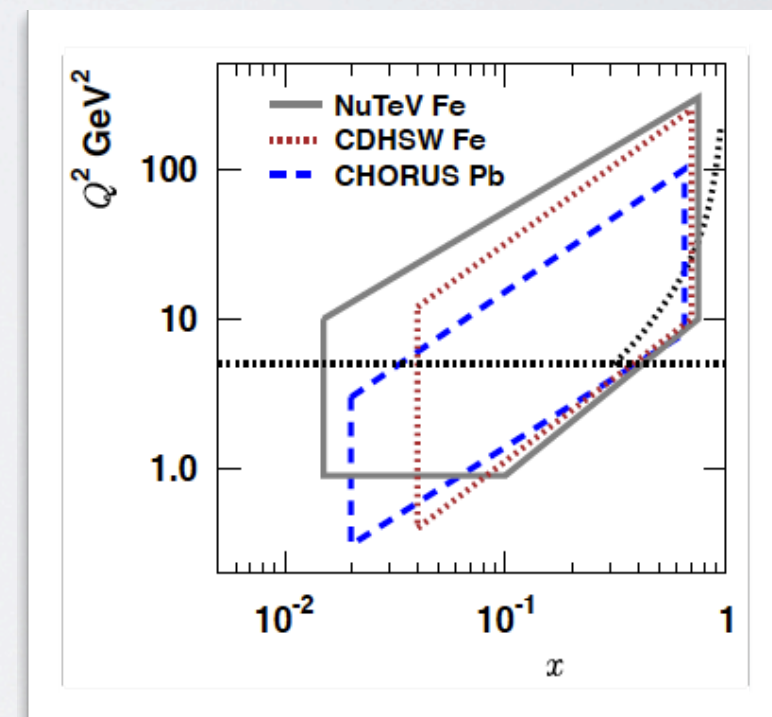
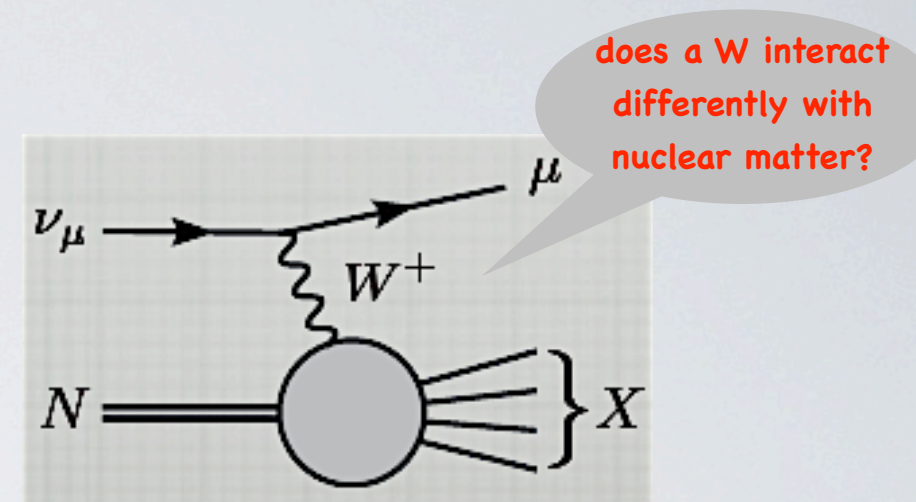
$$F_3^{\nu A}(x_N) \simeq [-(\bar{u}^A + \bar{c}^A) + d^A + s^A](x_N)$$

$$F_3^{\bar{\nu} A}(x_N) \simeq [u^A + c^A - (\bar{d}^A + \bar{s}^A)](x_N)$$

- experiments extract (under certain assumptions)

$$F_{2,3} \equiv (F_2^{\nu A} + F_2^{\bar{\nu} A})/2 \rightarrow \begin{cases} \bullet F_2 \text{ probes total quark singlet} \\ \bullet F_3 \text{ probes sum of valence PDFs} \end{cases}$$

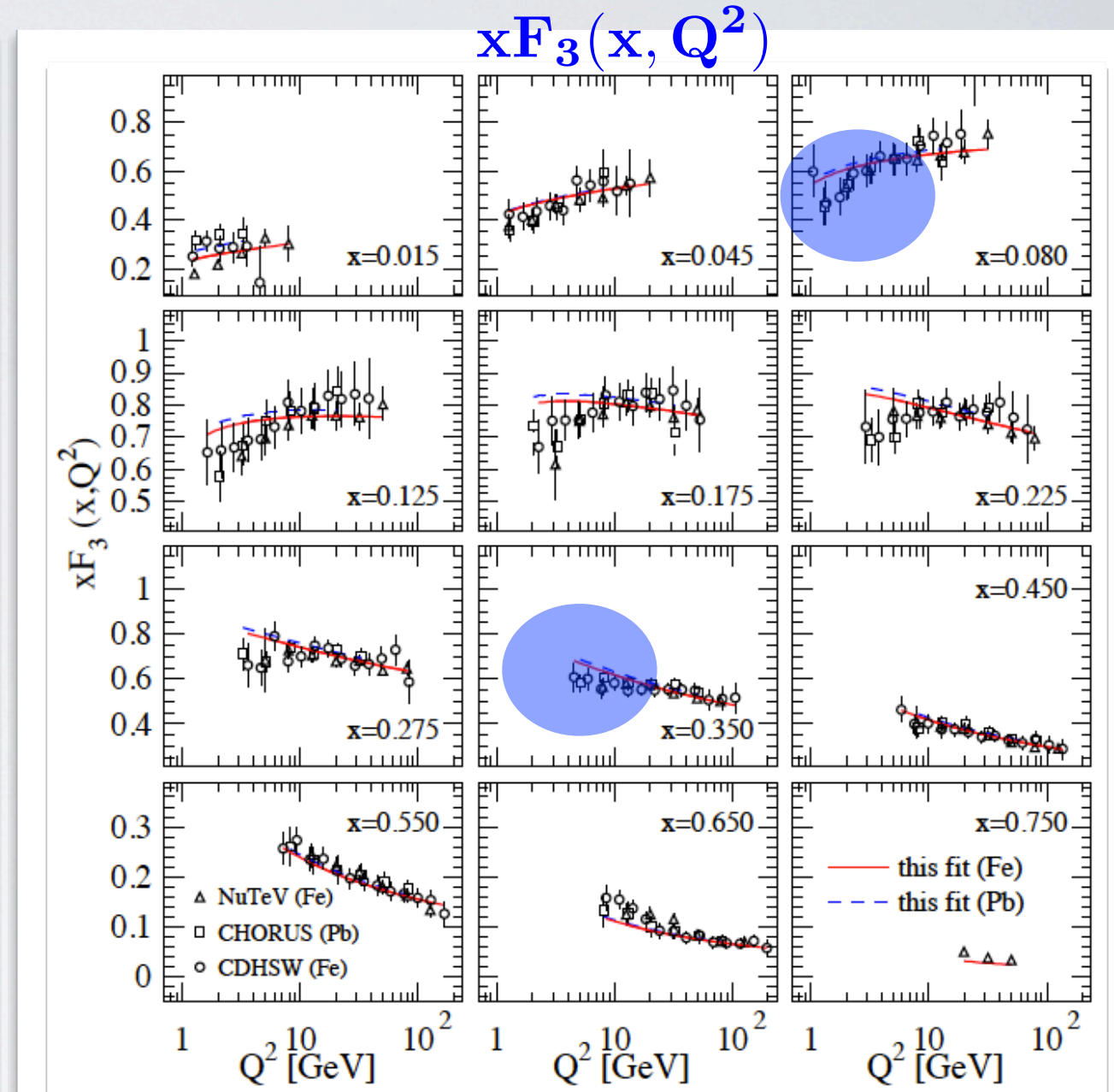
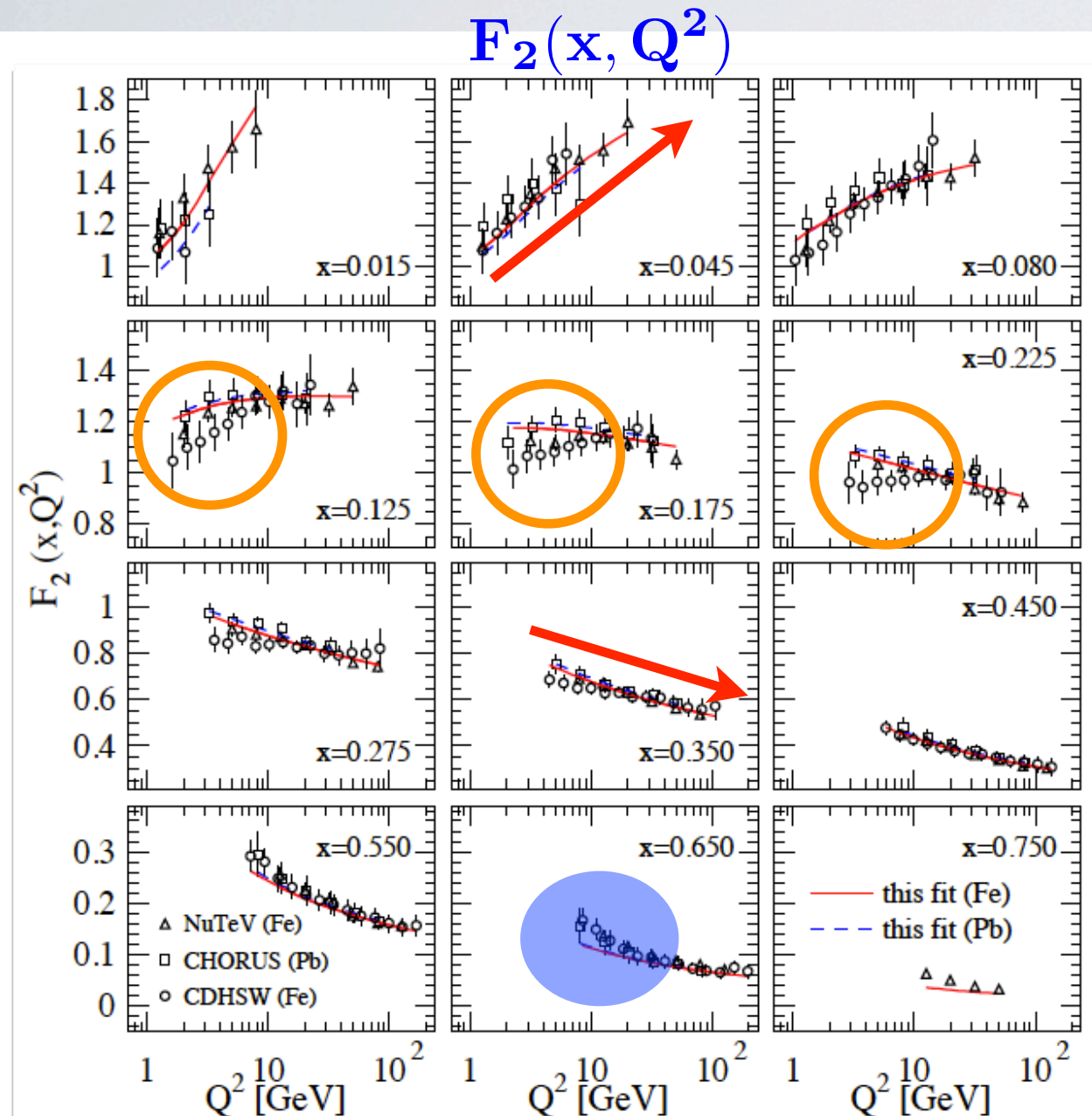
potential tension
with what we have
learned from NC DIS



kinematics overlaps with
charged lepton DIS data

DSSZ: cc neutrino DIS data

find: data remarkably well reproduced by fit $\chi^2 = 488.2/532\text{pts.}$



- ▶ absolute cross sections rather than ratios -> more sensitive to set of proton PDF in R_i^A (incl. as theor. uncertainty)
- ▶ data feature typical pattern of scaling violations
- ▶ slope of CDHSW data does not match with other data

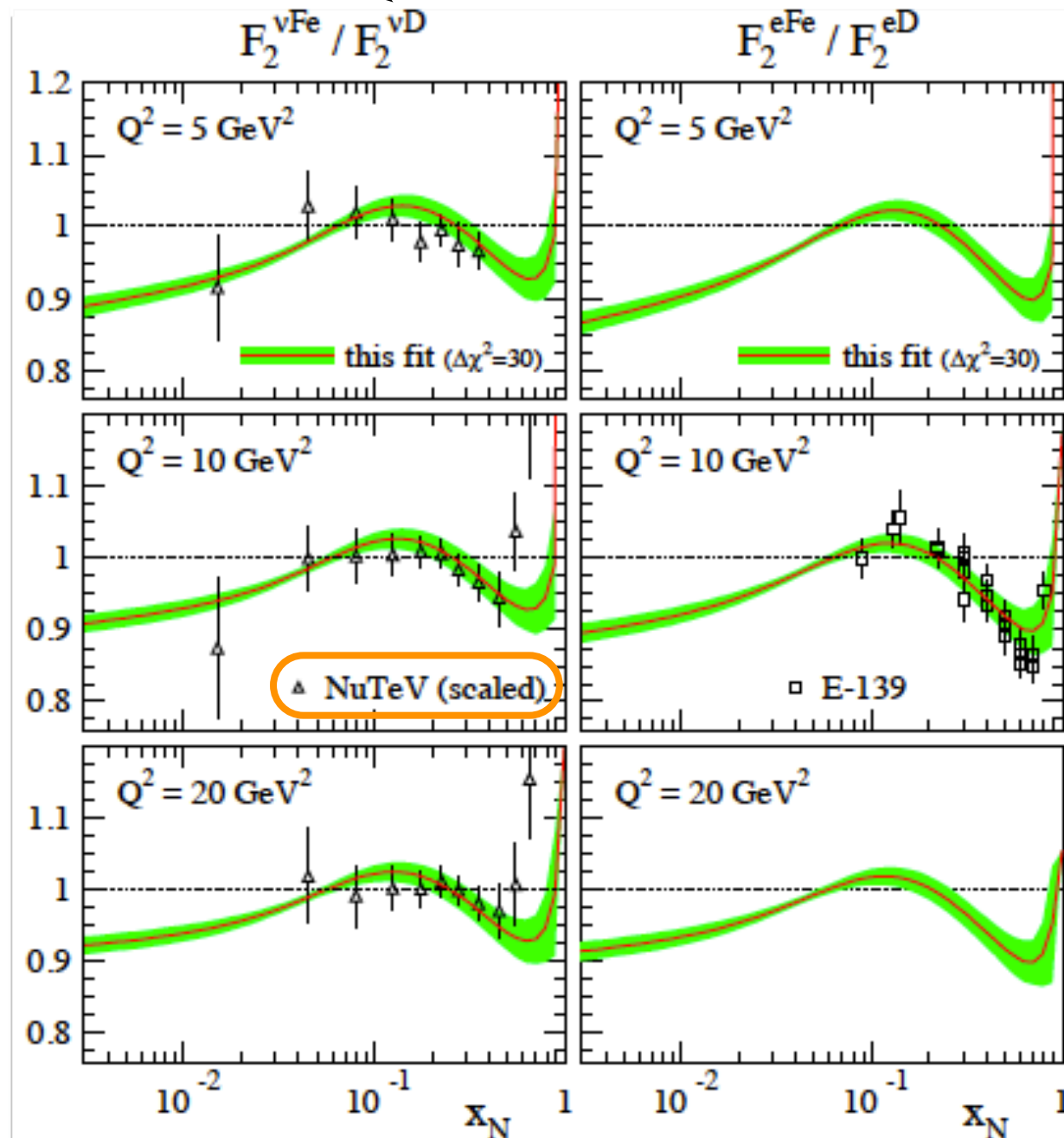
some mild tensions

often with CDHSW data

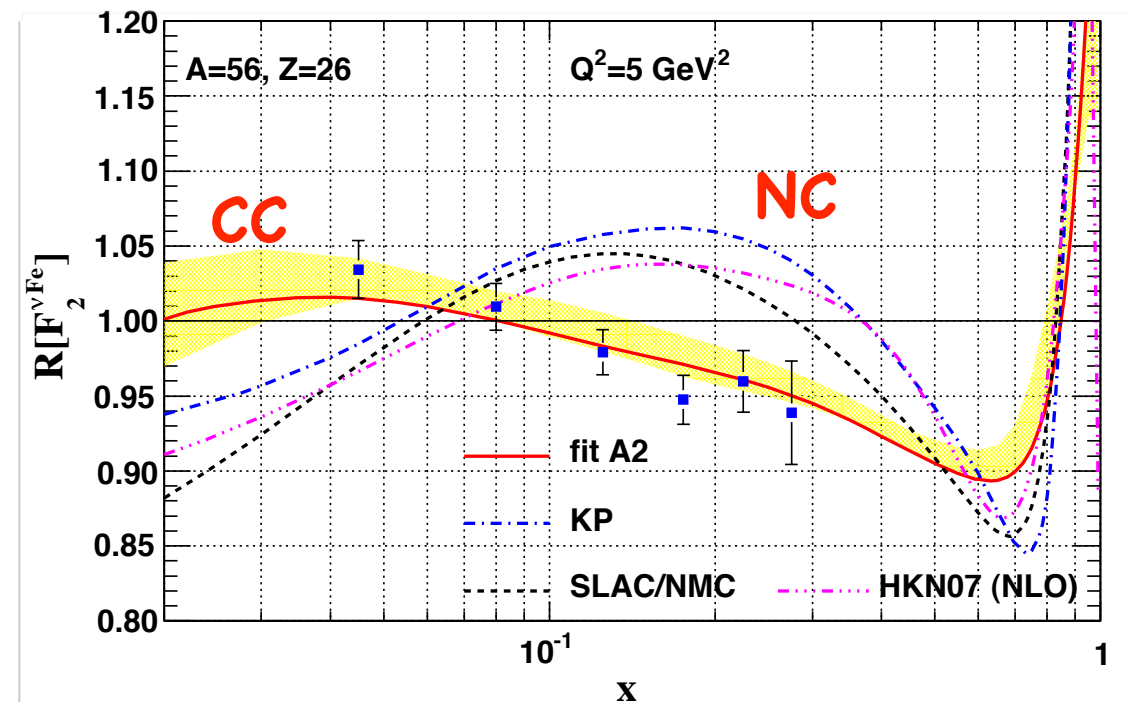
DSSZ: cc neutrino DIS data

no indication for factorization breaking

find same pattern of nuclear effects for CC and NC DIS



at variance with nCTEQ result



- “theoretical data”: $F_2^{\nu\text{D}}$ not measured
- nCTEQ fits to cross sections not str. fcts.
- also EPS finds compatible nuclear effects (no re-fit including CC DIS yet)

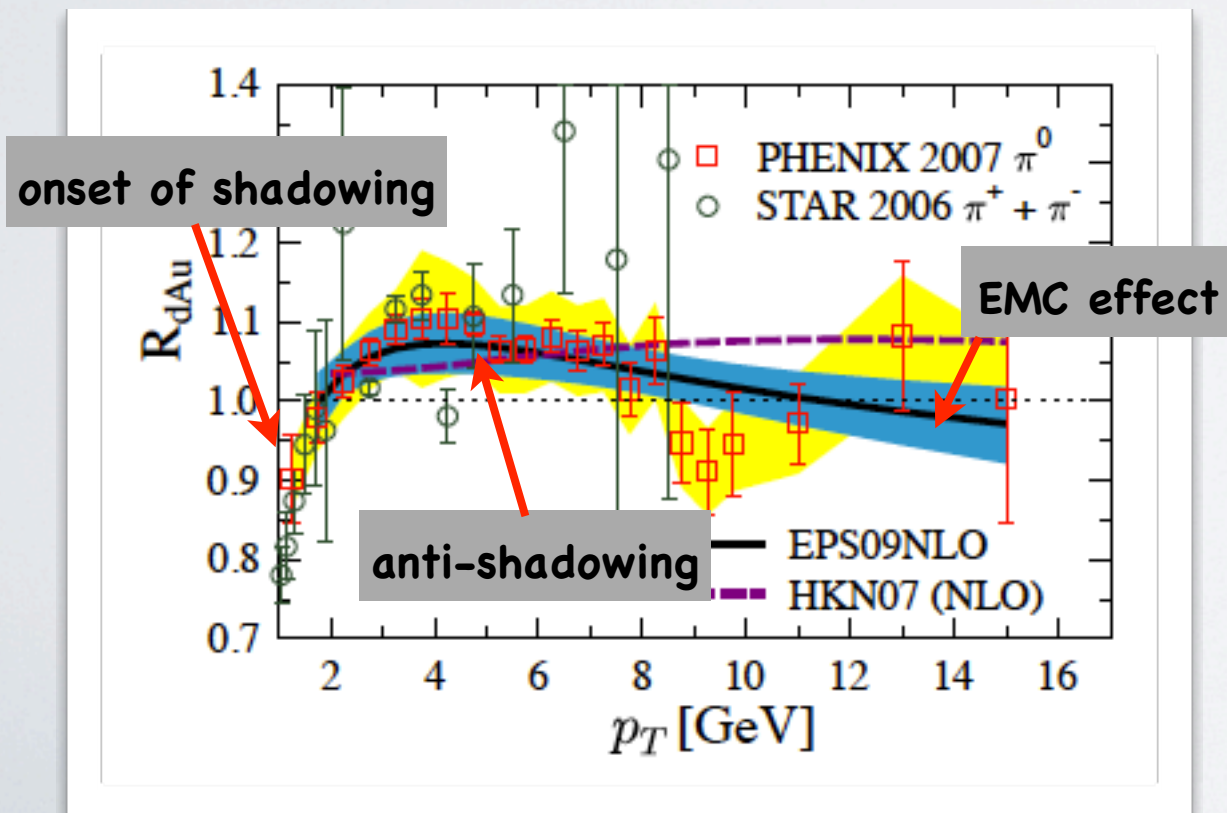
DSSZ: pion production dAu

most difficult probe to analyze (yet, perhaps one of the most interesting ones)

$$d\sigma_{dA \rightarrow \pi X}^A = \sum_{ijk} f_i^d \otimes f_j^A \otimes d\hat{\sigma}_{ij \rightarrow kX} \otimes D_k^{A,\pi}$$

free proton PDF "known" (points to f_i^d)
 wanted (points to f_j^A)
 known to NLO many contributing subprocesses (points to $d\hat{\sigma}_{ij \rightarrow kX}$)
 fragmentation functions fairly well known for pions (points to $D_k^{A,\pi}$)
 but what about possible nuclear modifications? can have an impact even if small (points to the entire equation with a large red question mark)

mid-rapidity neutral pion data from PHENIX and STAR first analyzed in EPS fit



- fit to min. bias ratio $R_{dAu}^\pi = \frac{\frac{1}{2A} d^2\sigma_{dAu}/dp_T dy}{d^2\sigma_{pp}/dp_T dy}$
- use up-to-date vacuum fragmentation functions
DSS: de Florian, R.S., Stratmann - include RHIC pp data
- find BIG impact on gluon nPDF

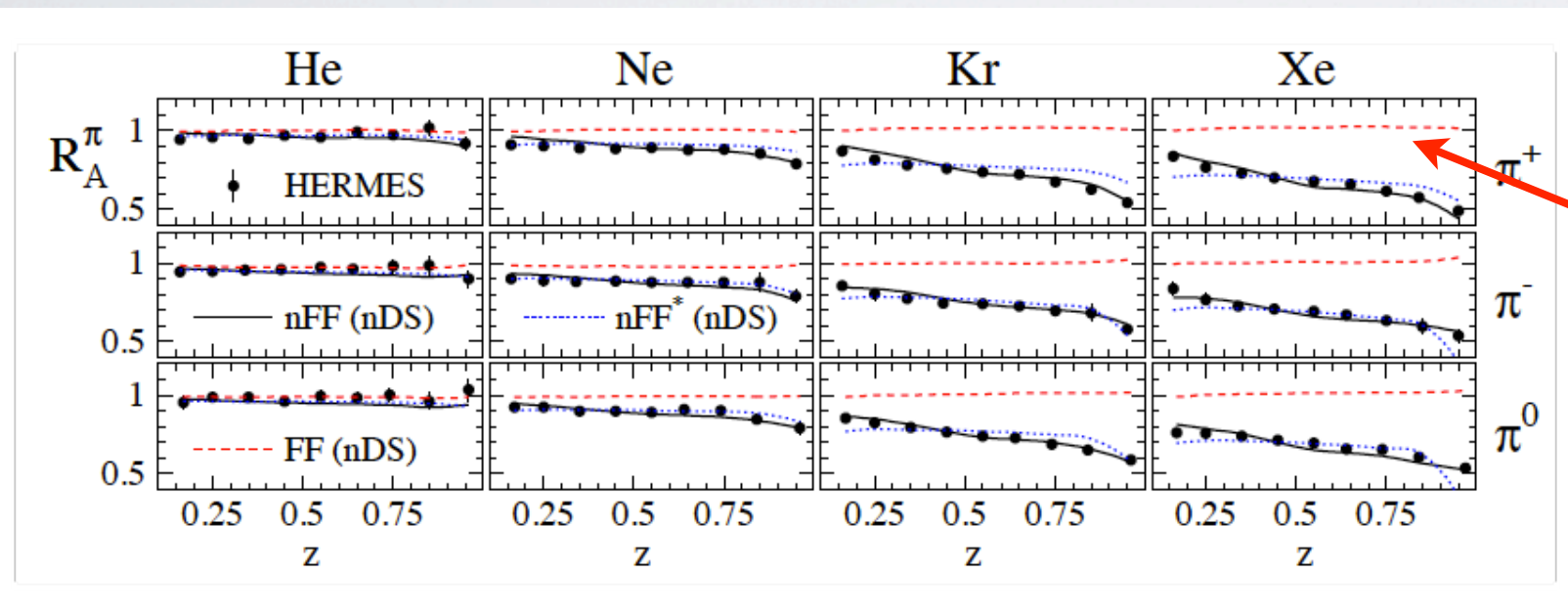
potential caveat: need to assign large weight to dAu data in fit

DSSZ: pion production dAu

what is different in DSSZ analysis

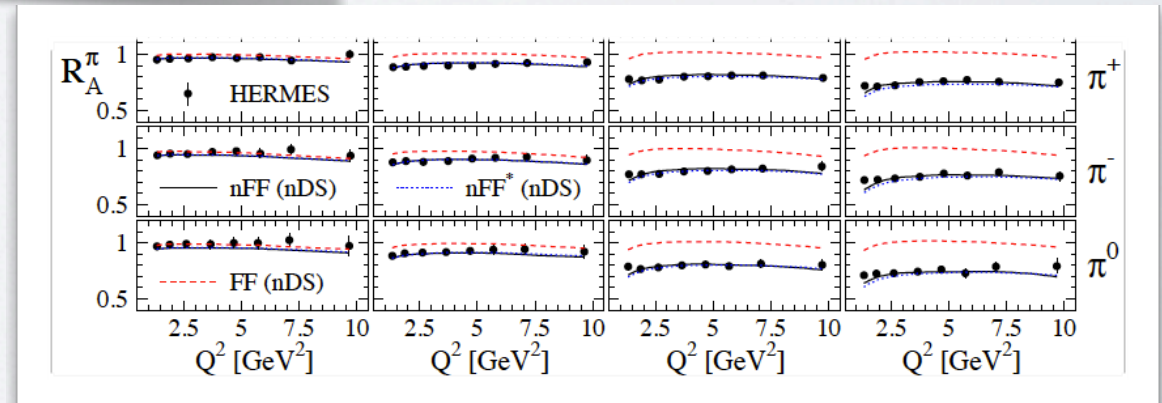
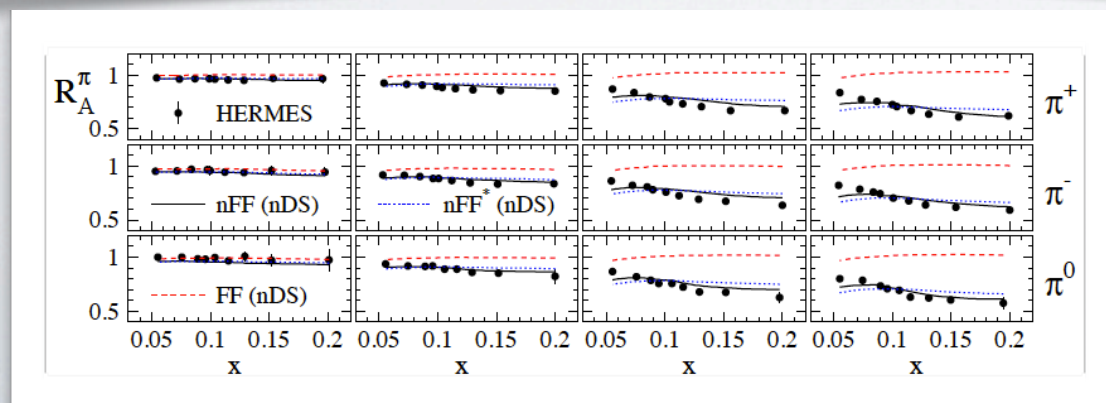
- ✓ more data, including also charged pions from STAR
- ✓ no artificially large weight w.r.t. other data sets
- ✓ try to estimate impact of modifications in hadronization

fragmentation in a medium – what is known ?



- effects known to be large in eA
- cannot be described as an initial-state effect (= nPDFs)
- hadron attenuation increases with A and z (rather flat in x and Q²)

HERMES



DSSZ: nFFs

how to model fragmentation in a medium ?

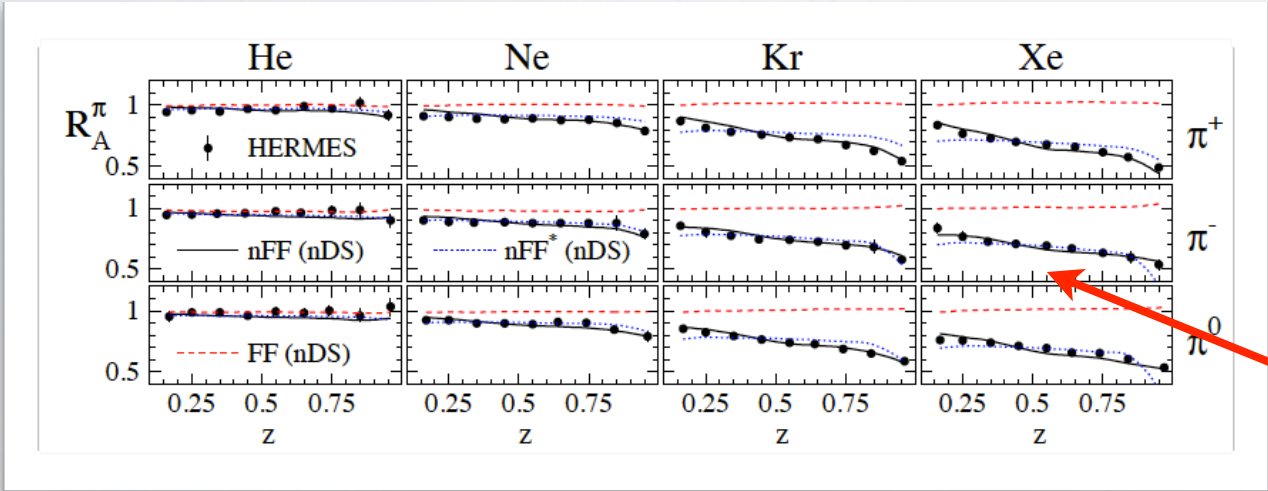
bold attempt: extend FFs to medium modified FFs ("in the background of a nucleus A") R.S, Stratmann, Zurita
0912.1311

choose convolution ansatz to modify vacuum FFs

DSS vacuum FFs

$$D_{i/A}^H(z, Q_0) = \int_z^1 \frac{dy}{y} W_i(y, A) D_i^H\left(\frac{z}{y}, Q_0\right)$$

from fit to HERMES and RHIC dAu pion data



works well

Experiment	A	H	Data type	Data points	χ^2
HERMES [6]	He,Ne,Kr,Xe	π^+	z	36	39.3
		π^-	z	36	23.0
		π^0	z	36	27.4
		π^+	x	36	69.4
		π^-	x	36	55.4
		π^0	x	36	49.7
		π^+	Q^2	32	21.0
		π^-	Q^2	32	27.1
		π^0	Q^2	32	34.7
PHENIX [14]	Au	π^0	p_T	22	13.7
STAR (prel.) [16]	Au	π^0	p_T	13	12.8
STAR [15]	Au	π^\pm	p_T	84	22.5
Total				381	396.0

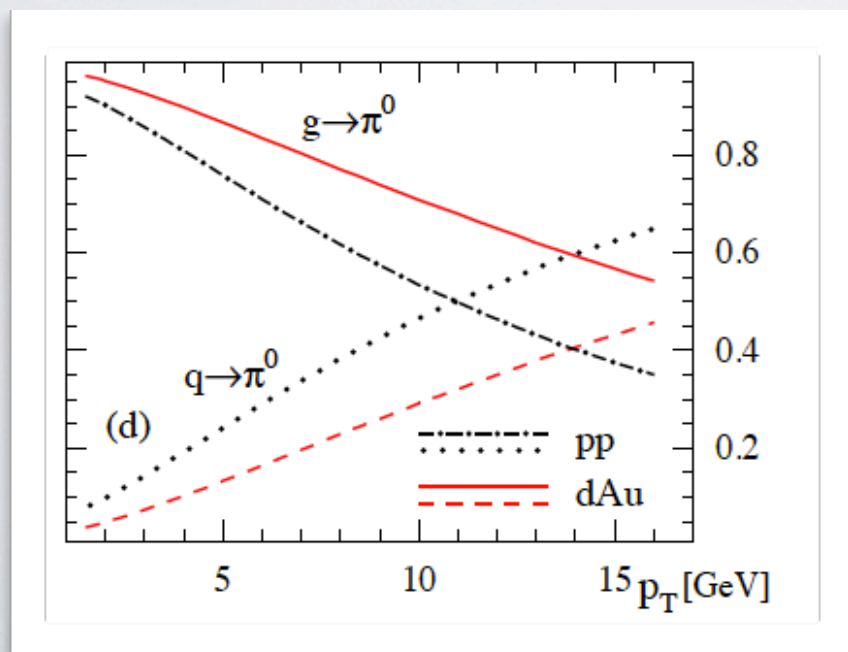
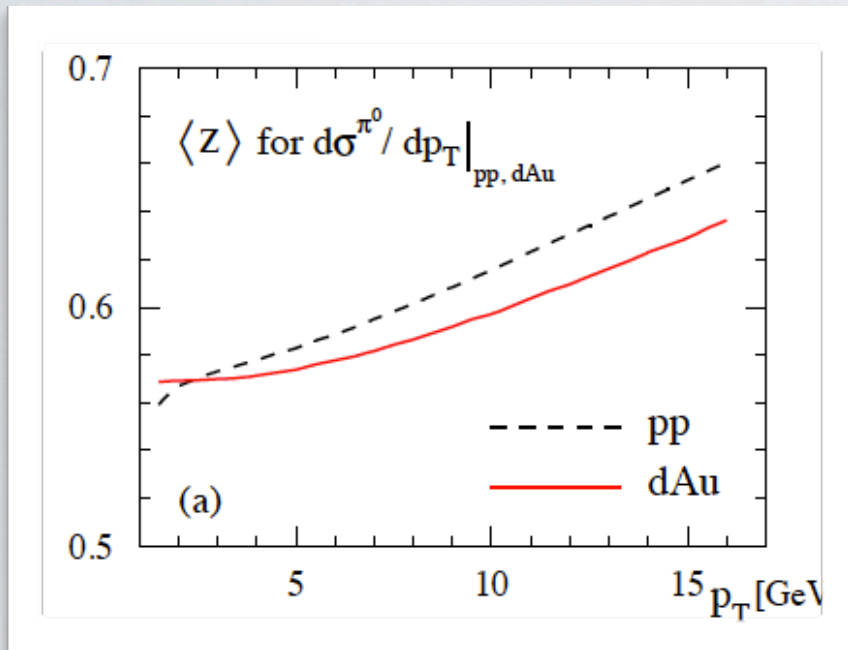
find:

- ▶ suppressed quark -> pion fragmentation (incr. with A)
- ▶ mildly enhanced gluon fragmentation around $z=0.5$

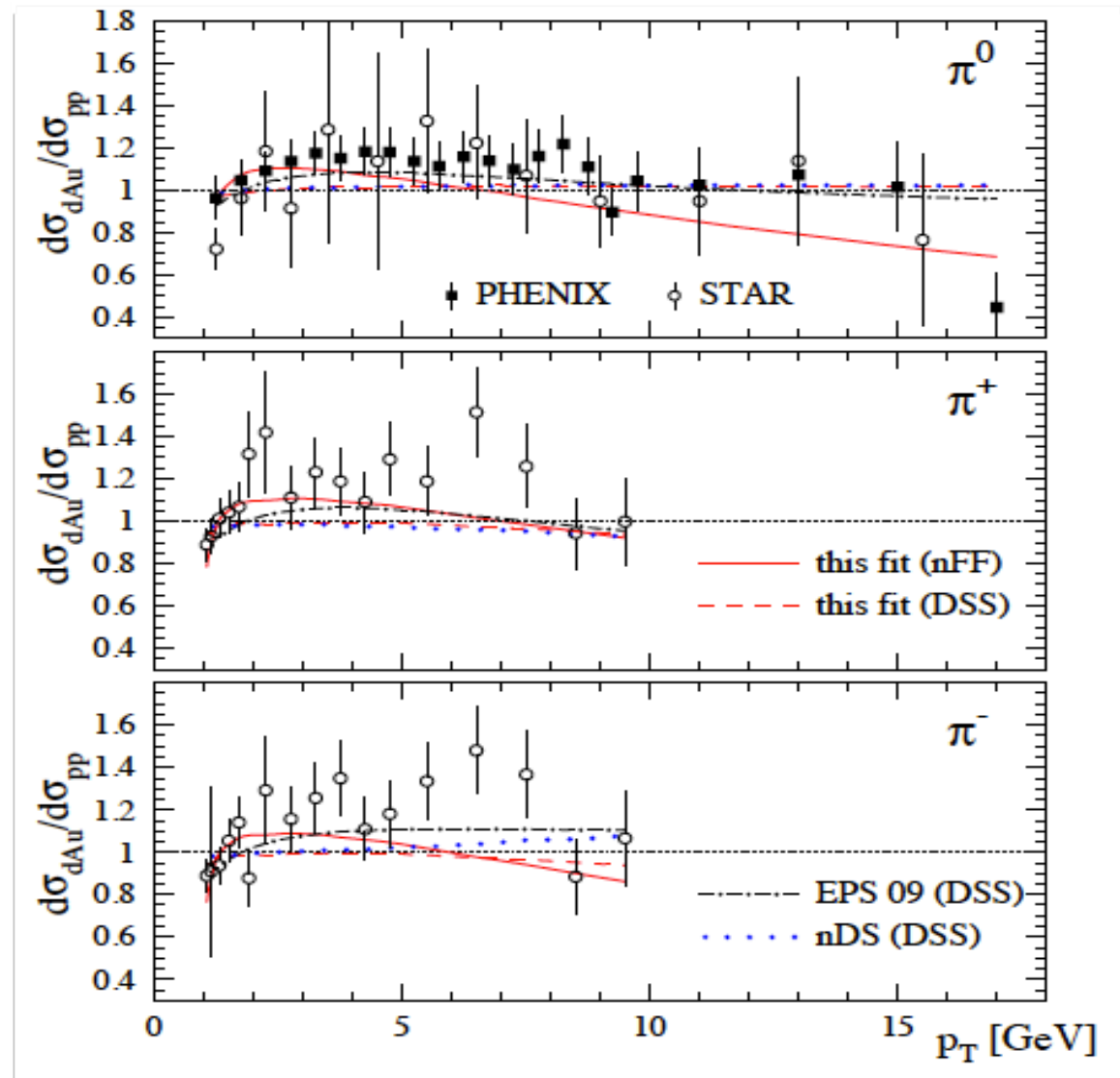
use both DSS vacuum and effective nuclear FFs in DSSZ nPDF analysis

DSSZ: mid rapidity (again)

at RHIC (mid rapidity) we probe large z
and mostly pions from gluons



result of our nPDF fit



► good fit within large exp. uncertainties

► choice of FF has some impact (but not too much)

$$\chi^2 : 68.3 \text{ (nFF)} \rightarrow 83.6 \text{ (DSS)}$$

► unlike EPS fit, limited impact on gluon (no weight factor)

DSSZ: forward rapidity

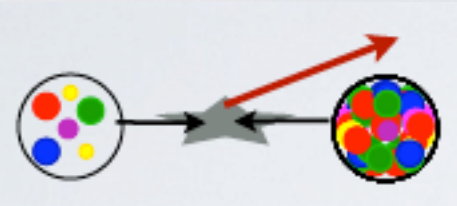
why interesting

- ▶ allows to access smaller x in nucleus
- ▶ gets one closer to the region where one expects saturation effects

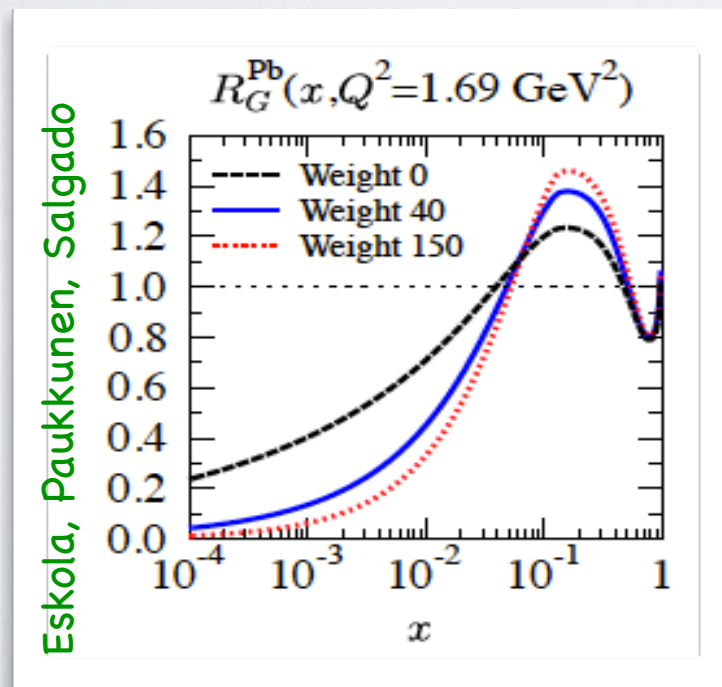
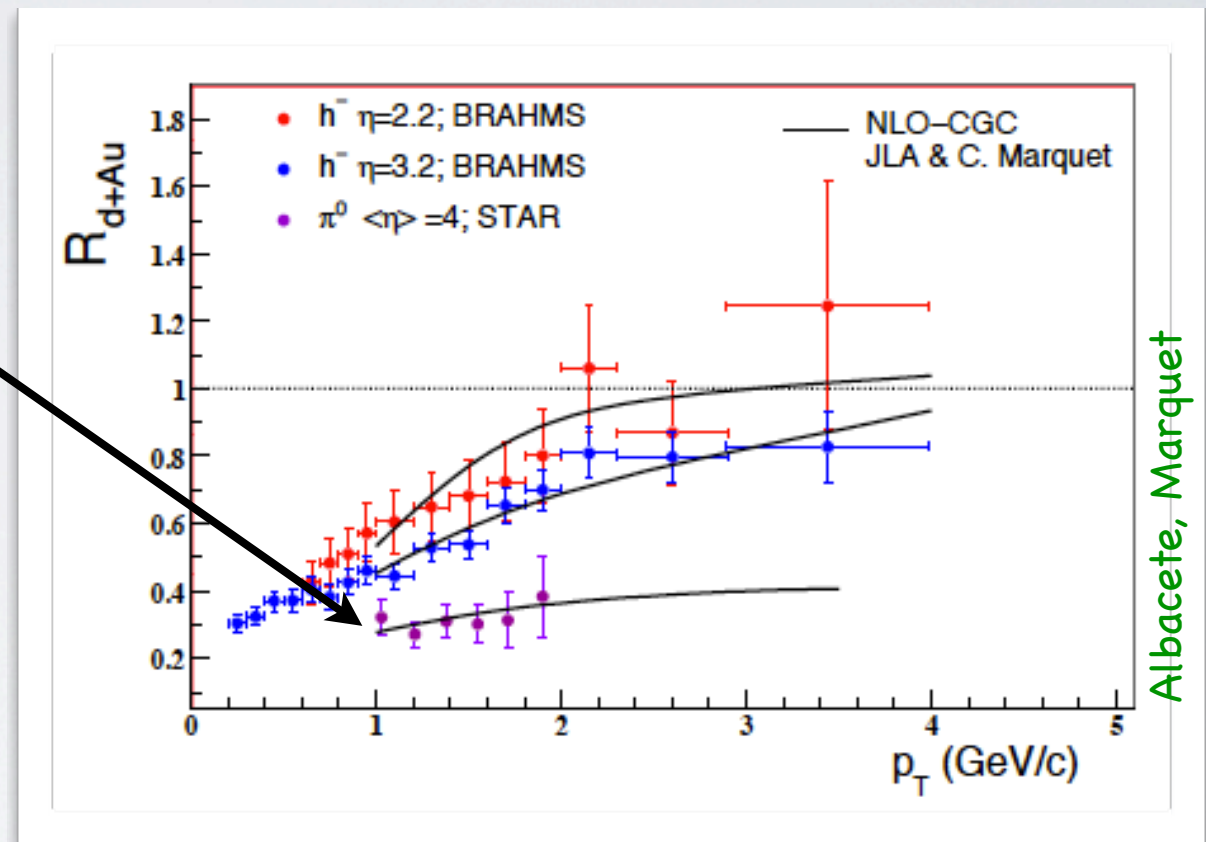
data indicate strong suppression of gluons at small x and low scales

forward suppression well described within **non-linear rcBK evolution (CGC)**

what does it take to describe it with nPDFs



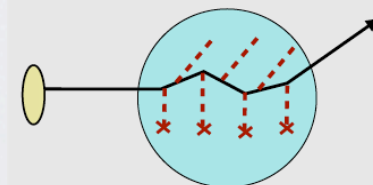
$$x_{1,2} \simeq \frac{p_T}{\sqrt{s}} e^{\pm y}$$



- ▶ need humongous shadowing at a scale of about 1 GeV



could be much less if final-state effects are relevant
advocated by Frankfurt, Strikman; Kopeliovich; ...



DSSZ:AA collisions *no, thanks*



many observables of interest involve

small p_T , global properties, centrality dependence,

- **nPDFs are collinear objects**

there is no impact parameter or other geometrical dependence

- **many observables in AA have no “hard scale”**

not amenable to pQCD calculations in standard factorizations

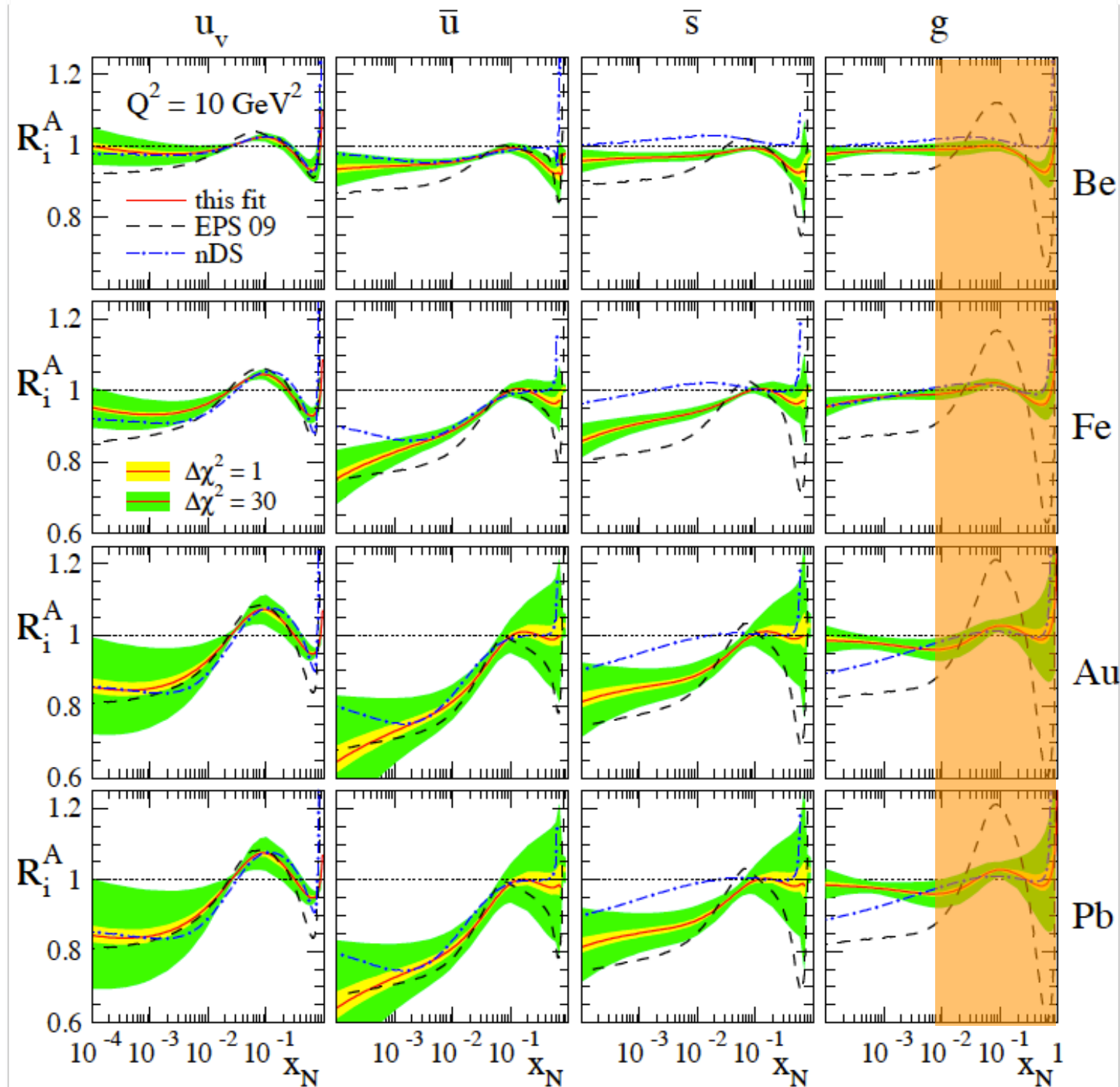
- **assuming factorization in AA is a stretch**

there might be some hard probes where things work out though

we do not touch AA data for the time being
nPDFs should be determined from probes in eA or pA
preferentially electromagnetic ones (free of hadronization issues)

DSSZ: nPDFs

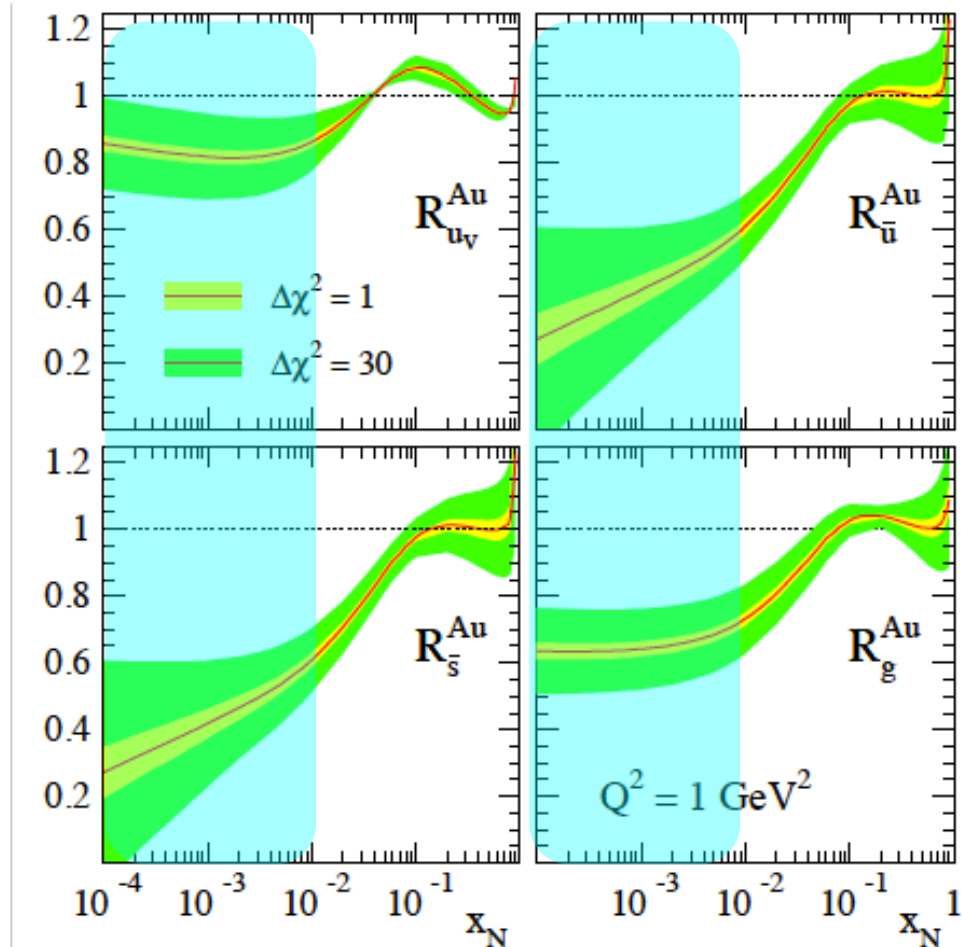
A dependence at $Q^2 = 10 \text{ GeV}^2$



- nuclear modifications increase with A
- good agreement with previous fits for $R_{u_v}^A$ and $R_{\bar{u}}^A$
- less so for $R_{\bar{s}}^A$ due to recent changes in free proton PDFs
- **MUCH less anti-shadowing and EMC effect than for EPS gluon** driven by the way dAu data are analyzed

DSSZ: nPDFs & uncertainties

uncertainties at input scale of 1 GeV (for gold nucleus)

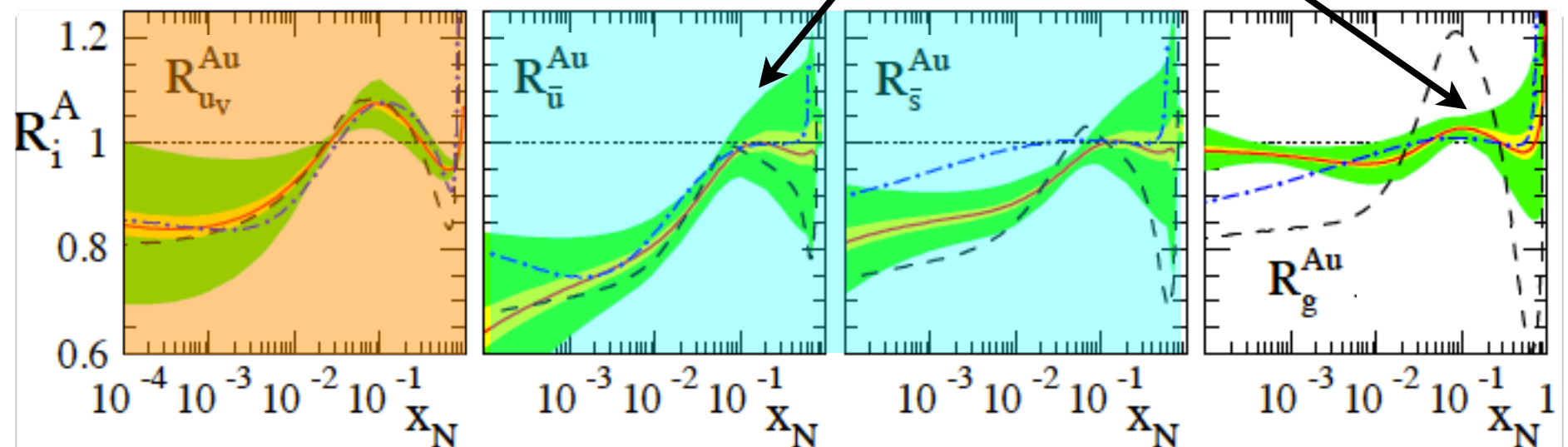


- uncertainties below 0.01 merely reflect extrapolation of chosen functional form
not constrained by any data

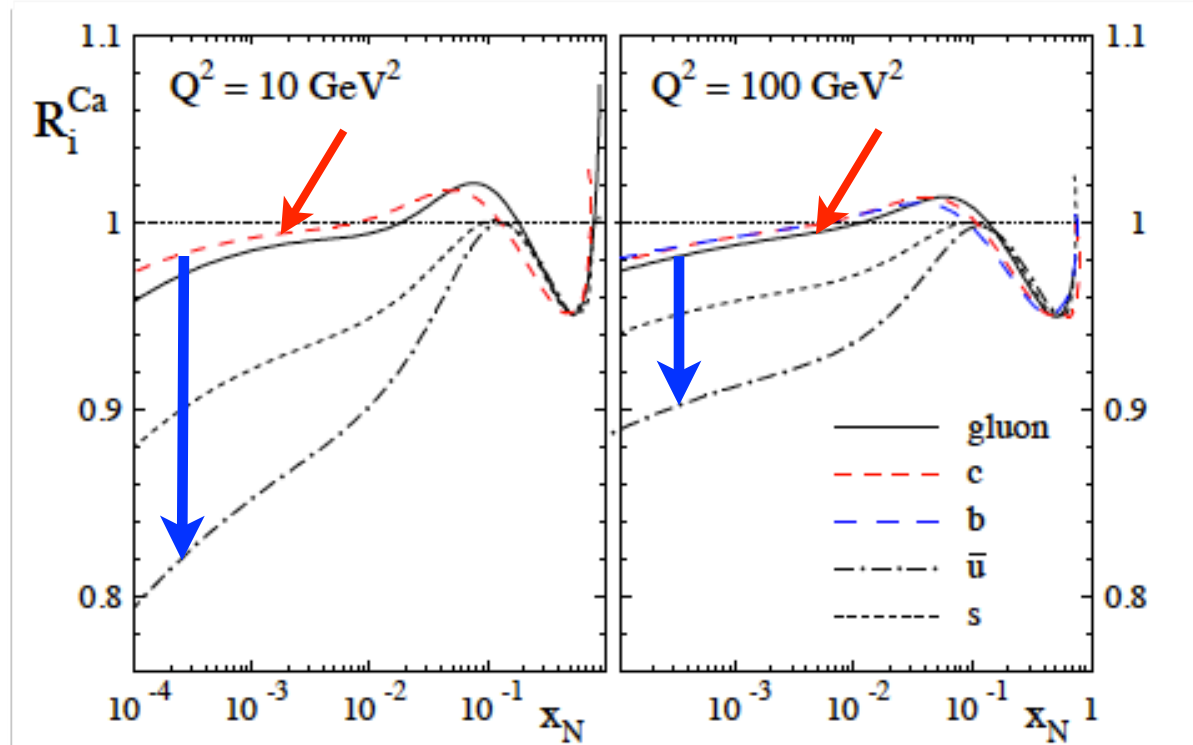


- nuclear modifications quickly diminish under evolution
- evolution imprints different nuclear effects on individual quark flavors
recall: we start with $R_u^A = R_d^A = R_s^A$
- $R_{u_v}^A$ exhibits textbook-like behavior
- little evidence for anti-shadowing in sea (and gluon)

evolve to 10 GeV^2



DSSZ: peculiarities



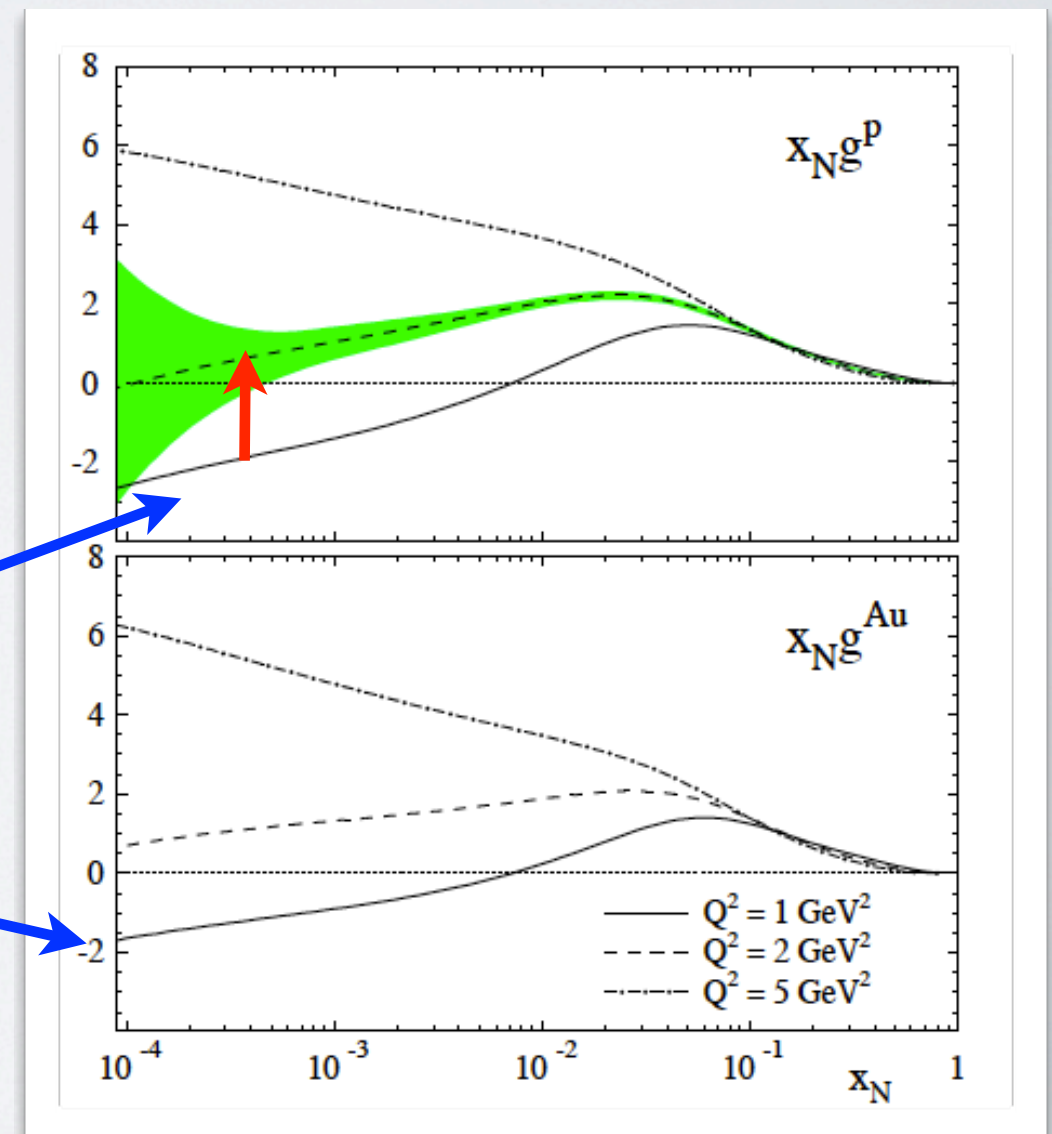
perturbatively generated charm and bottom nPDFs

- modifications for c,b follow closely the gluon no surprise, as they are generated from gluon splitting
- hierarchy in amount of low-x suppression: the stronger, the lighter the quark

the issue of “negative gluons”

- MSTW exercises the possibility of **negative gluons** at small x and low scales [improves their fit of HERA data]
not a problem since PDFs are not observables but F_L should stay positive
- evolution quickly pushes the gluon up
- our nPDF gluon is tied to the MSTW through R_g^A and gets negative too $\rightarrow R_g^A$ ill defined at low scales (nodes)

one must take trad. ratios R_i^A with a pinch of salt in NLO

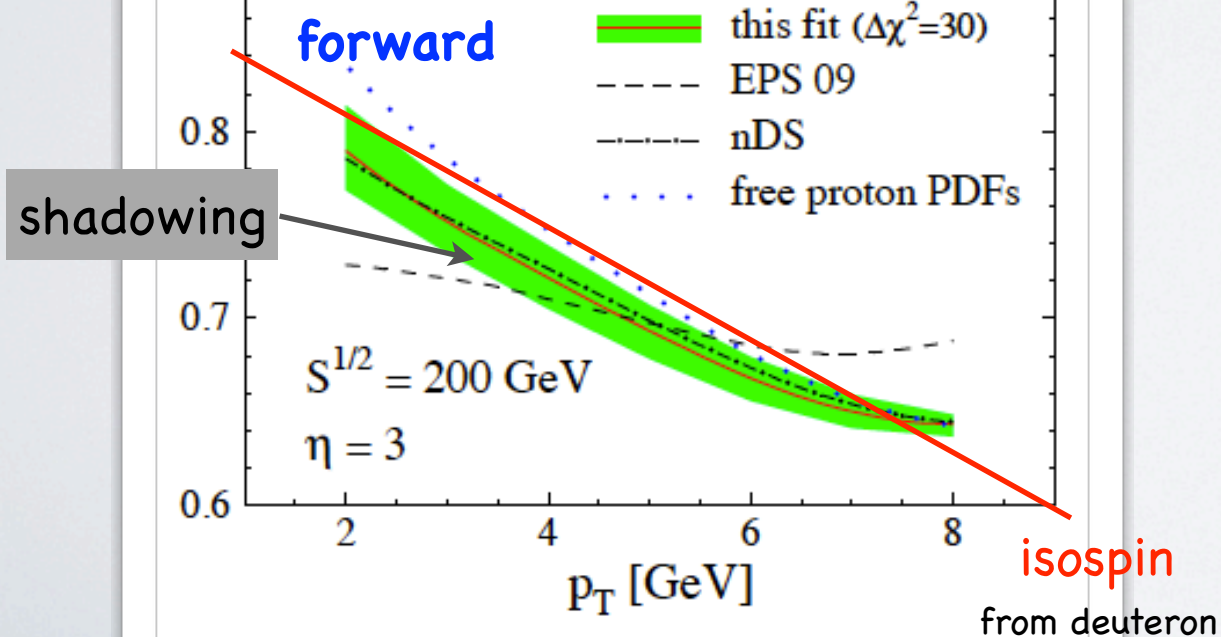
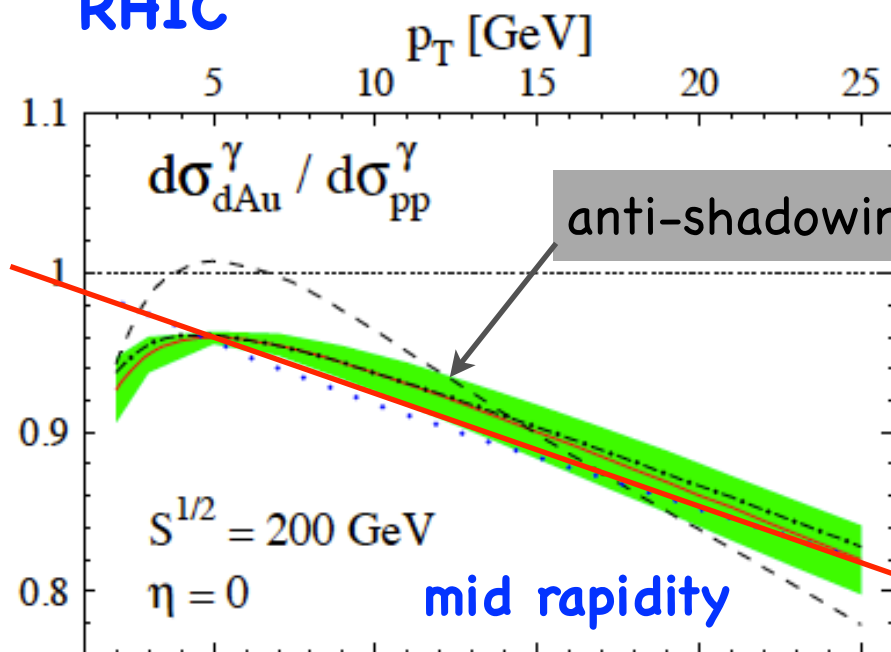


4.3 Future: prompt photons

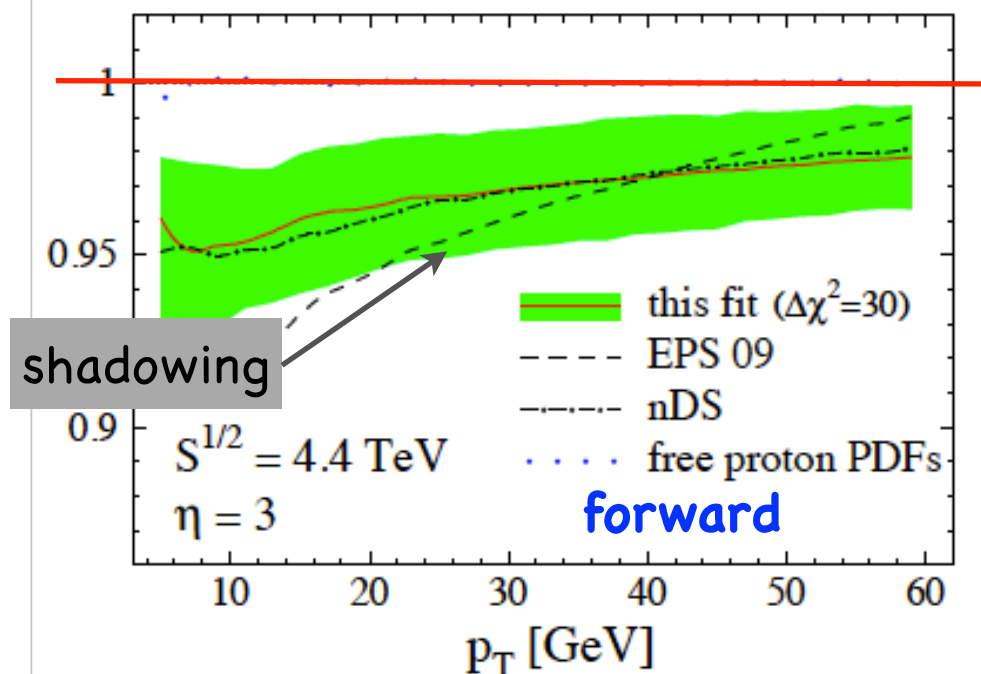
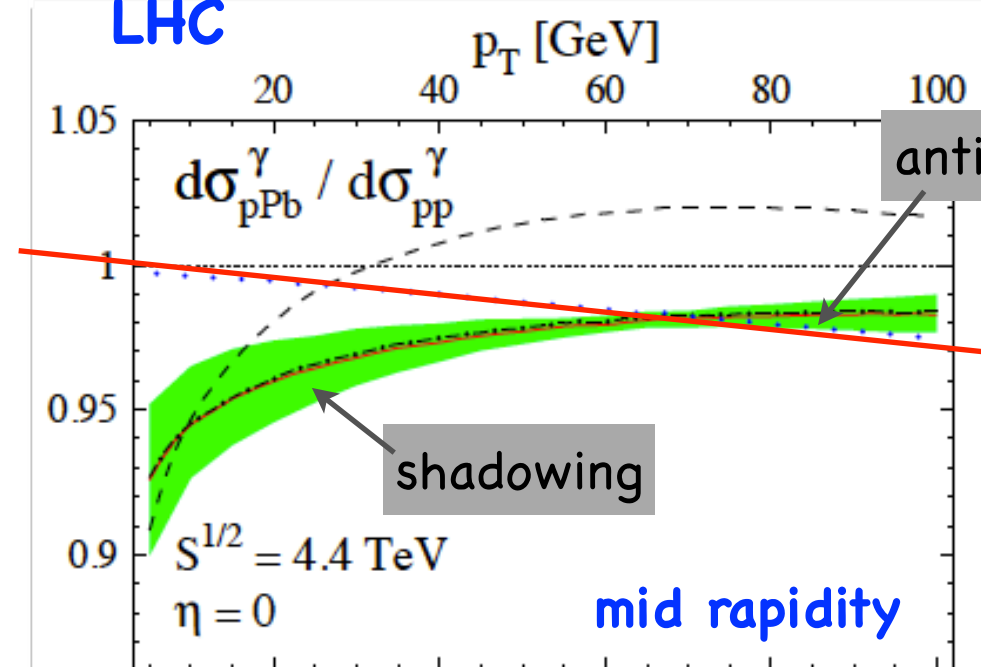
complication: "isospin effects" = dilution of u-quark density from neutrons $u^A(x) < u^P(x)$

→ ratio dAu/pp not unity even w/o nuclear modifications

RHIC



LHC



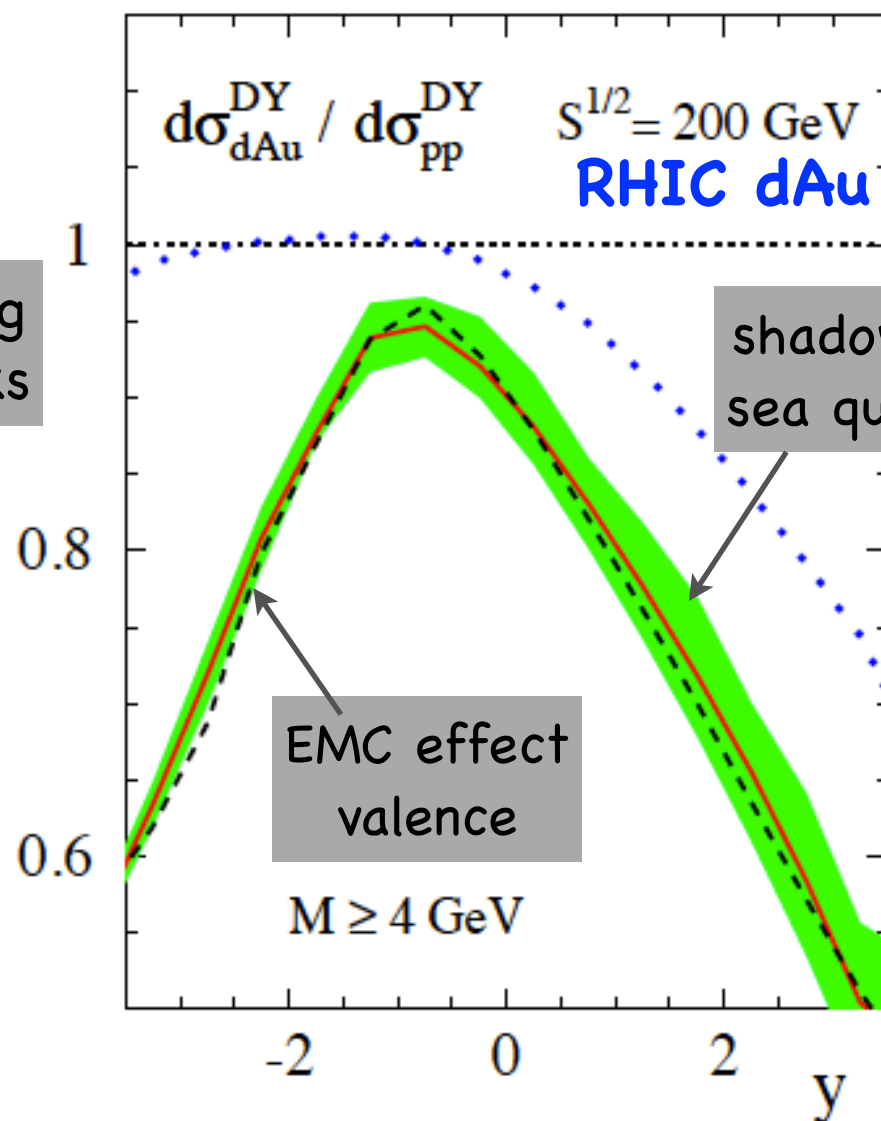
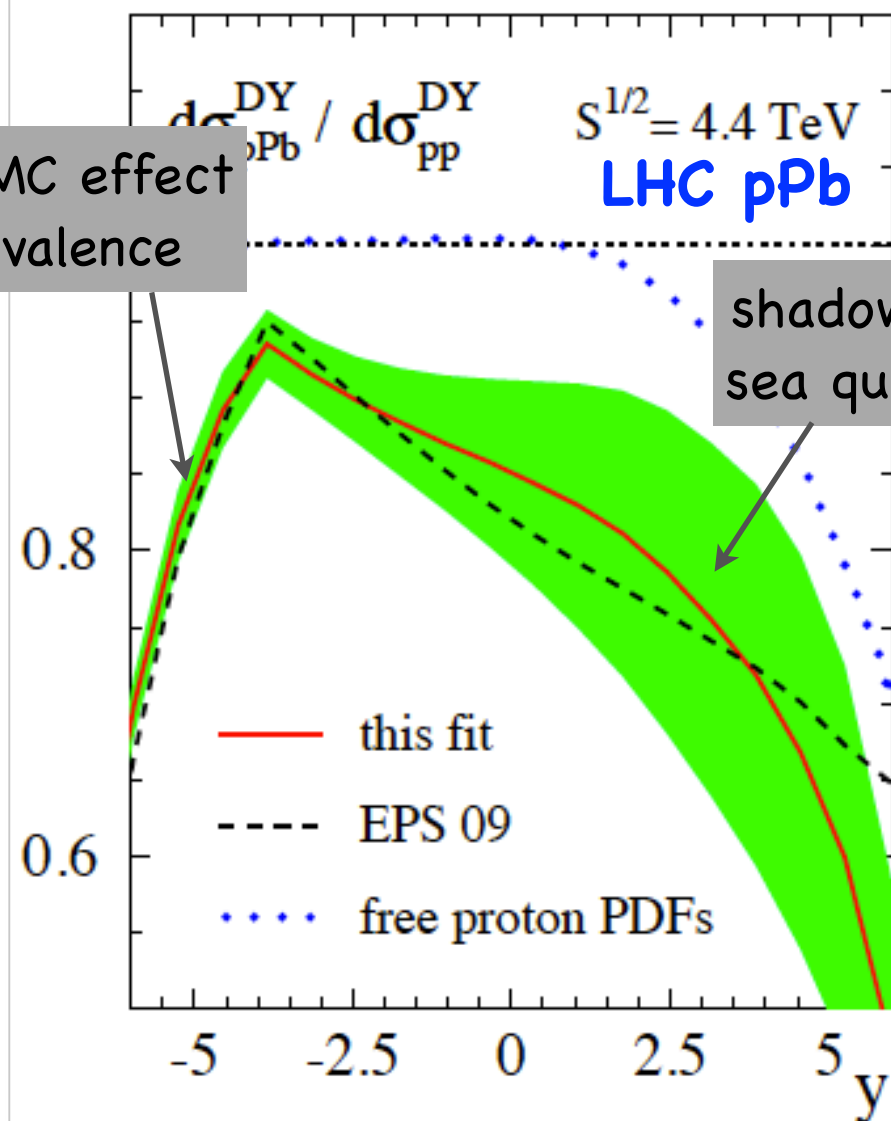
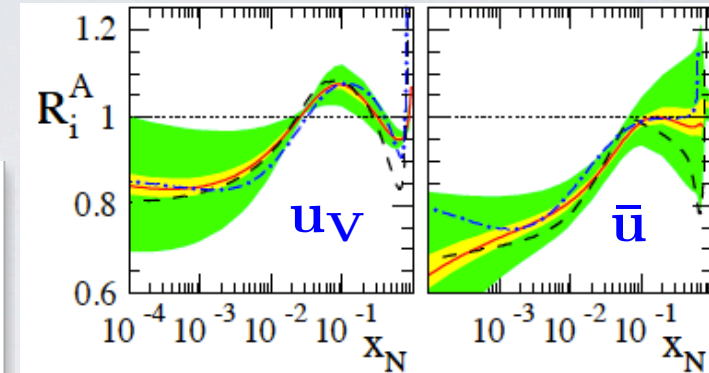
4.3 Future: Drell Yan lepton pairs

LO $d\sigma_{DY}^{pA} \propto e_u^2 [u(x_1)\bar{u}^A(x_2) + \bar{u}(x_1)u^A(x_2)]$
 $+ e_d^2 [d(x_1)\bar{d}^A(x_2) + \bar{d}(x_1)d^A(x_2)]$

large positive y

large negative y

$$x_{1,2} = \sqrt{M^2/s} e^{\pm y}$$



x reach at $y=3$

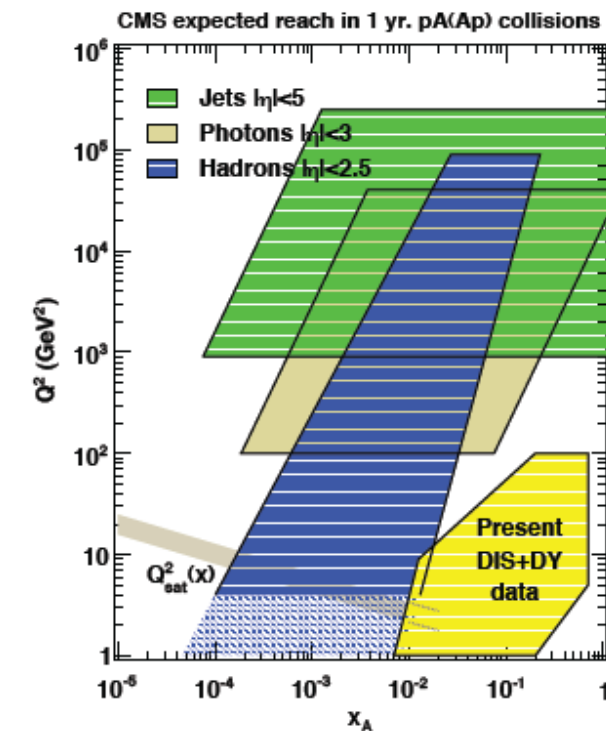
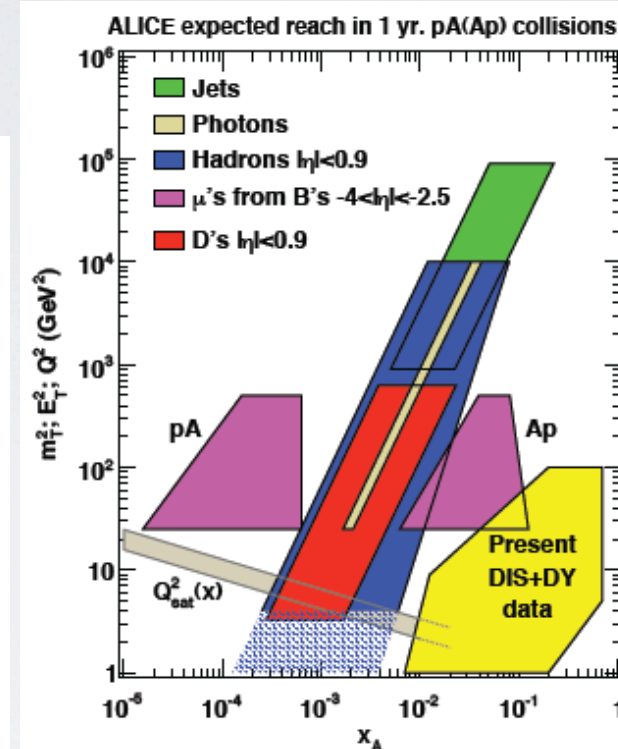
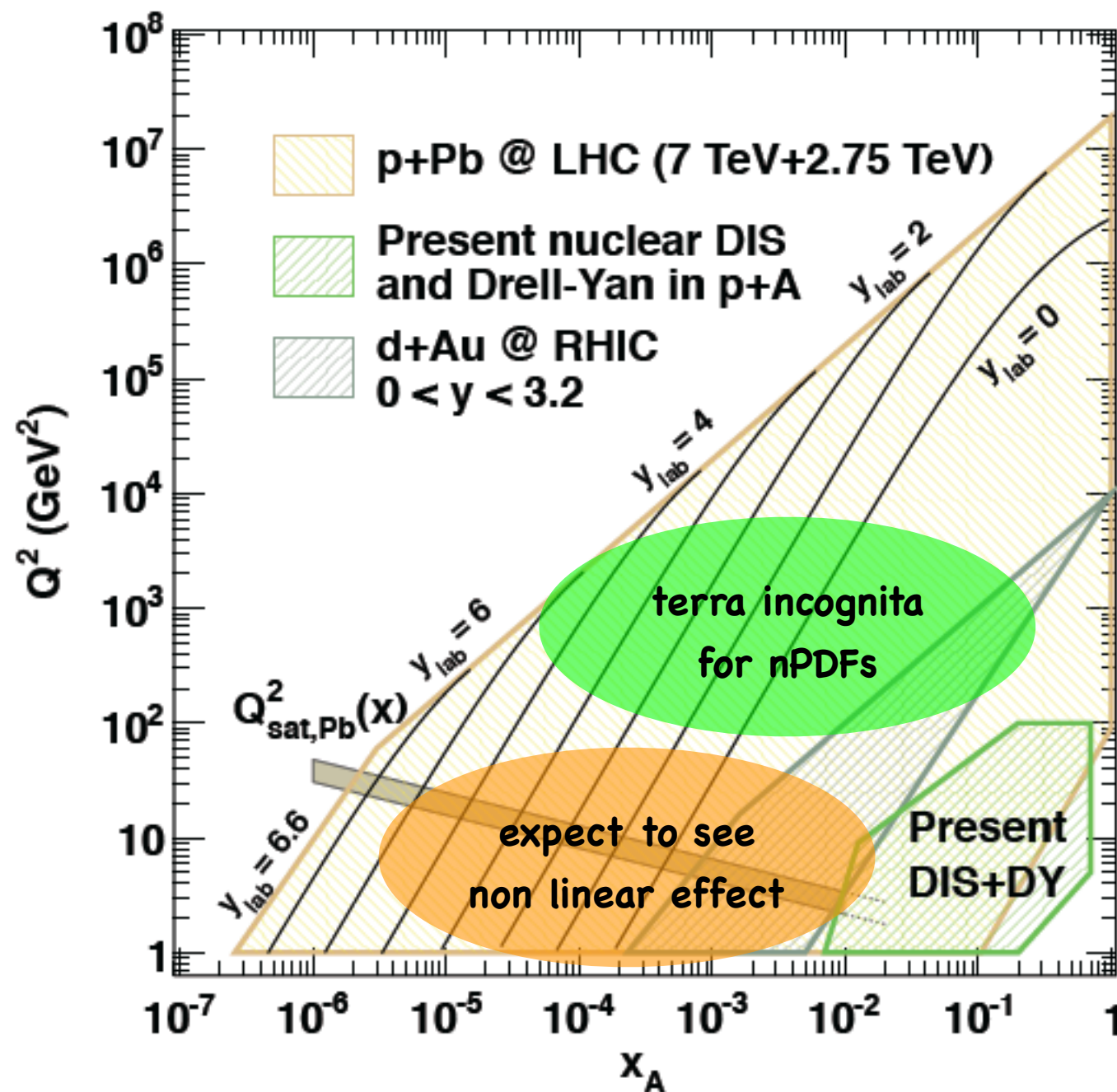
RHIC: $x_2 \simeq 10^{-3}$

LHC: $x_2 \simeq 5 \times 10^{-5}$

4.3 Future: pPb at LHC

see Salgado et al., 1105.3919

kinematic reach



- ▶ small x already accessible at mid rapidity
- ▶ many conceivable probes

expect great impact on nPDF fits

4.3 Future: eA at EIC & LHeC

PRECISION: direct access to nuclear partons through a leptonic probe

CONTROL: of the kinematic variables x , z , Q^2 over a very wide range

CLEANLINESS no fragments from another beam

HERA for nPDFs

Examples:

in addition to the standard low- x saturation, nuclear environment studies,

gluon nPDFs F_L scaling violations

high precision CC program to check factorization/universality of nPDFs

high precision program to check medium modified hadronization (nFFs)

...

1206.2913 1212.1701

Epilogue:

.... supposed to say something clever (bombastic?) about global analyses, PDFs, etc....

.... hmmm leave it as homework!

PDF customer satisfaction survey

How did you like global analyses?

What are their main pros and cons?

Do they affect your overall picture?

Can they improve your work? How?

Can your work improve them? How?

What would you like for Christmas?

Comments, complaints?

THANK YOU!

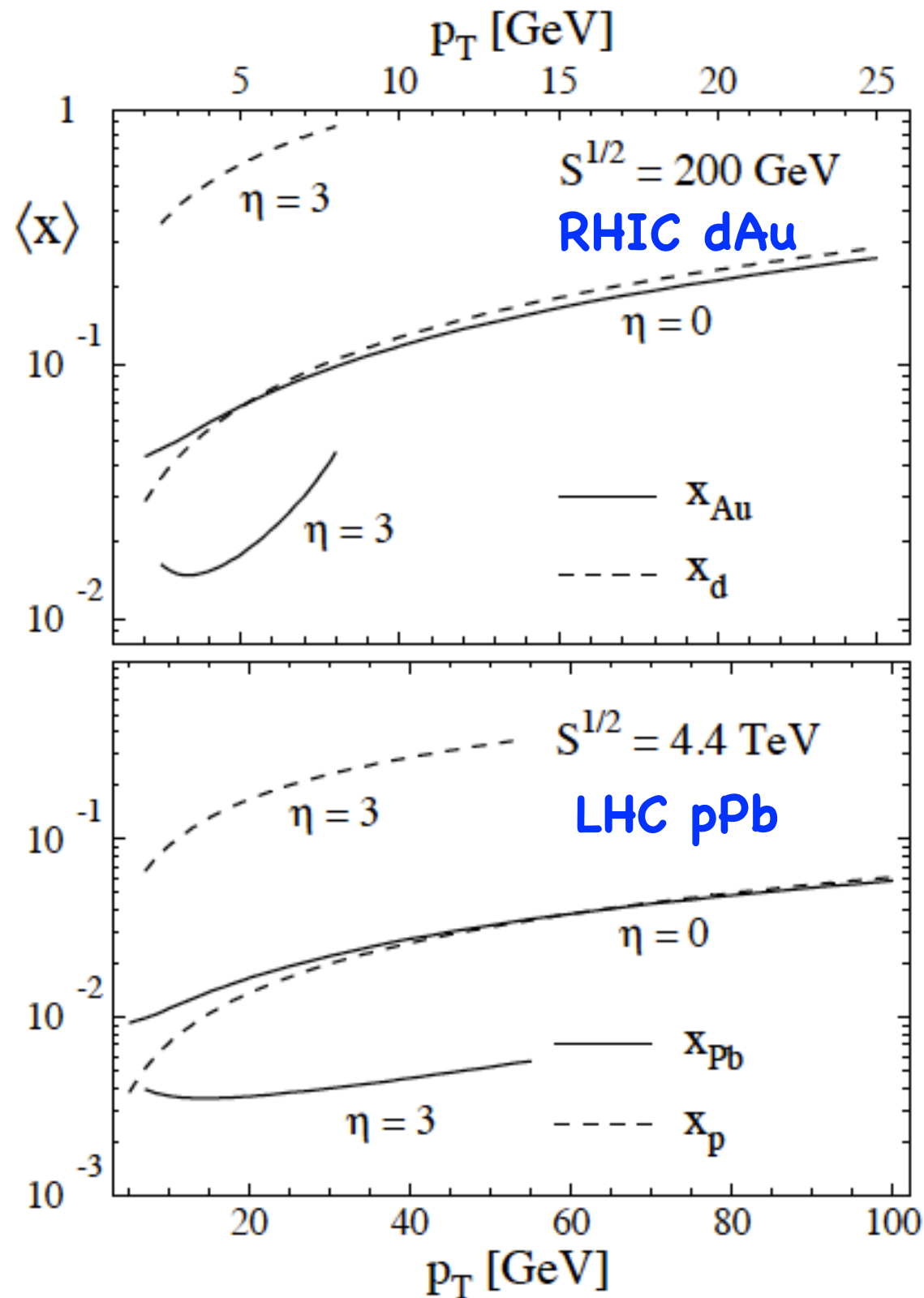
attention (patience!)

hospitality

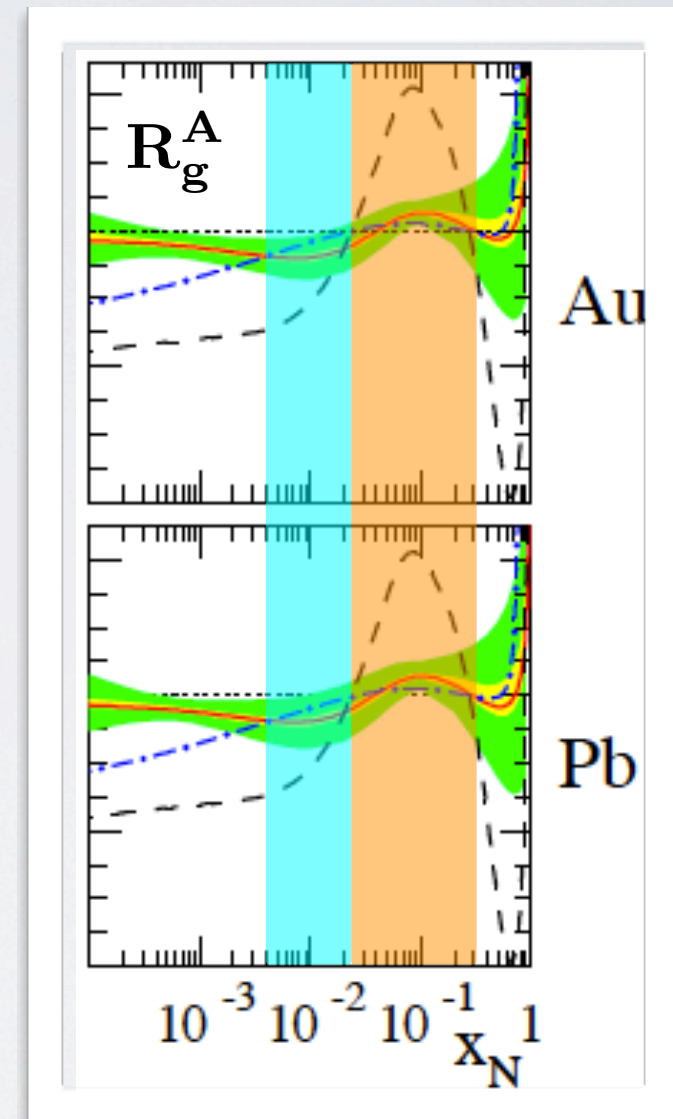
data!!

➡ sassot@df.uba.ar

4.3 Future: prompt photons



- can resolve characteristic differences between EPS and DSSZ gluons in anti-shadowing [and EMC] region



- can probe into shadowing region